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**LUNAR PHYSICAL PARAMETERS STUDY**

**PARTIAL REPORT NO. 12**

**MEASUREMENTS RELATED TO THERMAL CONDUCTIVITY**

**WORK PERFORMED UNDER J.P.L. CONTRACT NO. N-33552**

*TEXACO, INC.*

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MEASUREMENTS RELATED TO THERMAL CONDUCTIVITY

In support of development work done on a method for measuring the thermal diffusivity of the lunar surface, a mathematical study has been undertaken. This study was designed to verify the feasibility and design calculations previously reported<sup>1</sup>, and to provide some experience in calculation methods which might be used in the interpretation of data actually obtained.

The apparatus built for breadboard tests was used as the basis for calculation. The apparatus consisted of a flat circular disc, 12 in. in diameter, which was shielded from direct rays of the sun by a reflector directly above it. This allowed the disc to assume a lower equilibrium temperature than if it were in direct sunlight. The entire bottom surface of the disc was covered with an electrical heating element cut in a spiral of archimedes configuration. The lower surface was blackened to make the emittance as high as practical. A 1 in. diameter hole through the center of the disc allowed a radiometer mounted between the disc and reflector to view the area of the surface directly beneath the disc.

The calculations done were based on experiments which would be performed in the following manner:

1. An experiment performed during the lunar day - The disc, at a temperature lower than that of the surface, is placed

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<sup>1</sup>Lunar Physical Parameters Study, Partial Report No. 8, Design Calculations, Measurement of Thermal Diffusivity, Texaco Inc., May 8, 1961

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near the surface and parallel to it. The temperature of the disc is maintained constant<sup>2</sup> and the temperature of the surface is measured by the radiometer at a number of times, up to perhaps an hour after placement of the disc.

2. An experiment performed during the lunar night -

The disc is heated electrically and held at a constant uniform temperature. The disc is then positioned over the surface and the temperature measured by the radiometer at a number of times.

The actual calculations were approached in two phases. The first phase consisted of attempting to find reasonable sets of physical parameters which would allow calculation of temperatures and temperature gradients in the near-surface region. Results of such calculations could then be used as initial conditions for the calculation of the expected experimental results. These results also provide estimates of optimum times for performing the experiments, when the temperatures and temperature gradients are slowly varying with time over the surface area of interest. The second phase involved calculations of expected results of the experiments described earlier with various sets of assumed physical parameters.

Phase 1

The first phase calculations were based on a layered model.

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<sup>2</sup> Its temperature actually rises slightly, but with large heat capacity of the disc this is minimized and was ignored in the calculation. Some additional advantage might be obtained by supplying a small amount of thermostatted power to the heating element sufficient to balance radiation losses from the disc which can be cut off when the disc is positioned.

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The surface was assumed to be flat and of uniform character. The medium was treated as a semi-infinite slab made up of a finite number of parallel homogeneous, isotropic layers. Thus, the heat equation in the interior of any layer is;

$$\frac{\partial v(x,t)}{\partial t} = x \frac{\partial^2 v(x,t)}{\partial x^2} \quad (1)$$

where,

$x$  is the thermal diffusivity of the layer

$v$  is the temperature

$t$  is the time, and

$x$  is the depth variable.

At the surface,  $v(0,t)$  is known from astronomical data and the net heat flux conducted into the surface is

$$f(0,t) = \alpha_v q - \epsilon_s \sigma [v(0,t)]^4 \quad (2)$$

$\alpha_v$  = average reflectivity of surface in solar spectrum

$q$  = incident solar radiation of the form  $q_0 \sin \Omega t$  during the day and zero during the night

$q_0$  = solar constant

$\Omega$  = synodic frequency  $\times 2\pi$

$\epsilon_s$  = emittance of the surface in the infrared

$\sigma$  = Stefan-Boltzmann constant.

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The heat equation rewritten in terms of the flux becomes,

$$\rho c \frac{\partial v(x,t)}{\partial t} = - \frac{\partial f(x,t)}{\partial x} \quad (3a)$$

$$f(x,t) = -K \frac{\partial v(x,t)}{\partial x} \quad (3b)$$

$\rho$  = density of medium

$c$  = specific heat

$K$  = thermal conductivity =  $\rho pc$ .

A closed form solution of (3) is obtained through representation of the boundary conditions by the Fourier series

$$v(0,t) = a + \sum_{n=1}^{\infty} [(2p_n) \cos \omega_n t + (2q_n) \sin \omega_n t] \quad (4a)$$

$$f(0,t) = -K[b + \sum_{n=1}^{\infty} [(4r_n) \zeta_n \cos \omega_n t + (4s_n) \zeta_n \sin \omega_n t]] \quad (4b)$$

where,

$$\omega_n = n\Omega \text{ and } \zeta_n^2 = \frac{1}{2} \omega_n / \kappa.$$

We approximated this representation by a least square fit of the boundary data to Fourier sums containing a constant term and the first 20 sine-cosine pairs. Each of the sine-cosine pairs and the constant term can be considered separately. Then the results of these 21 solutions combined by superposition to give the desired solution to (3). In Appendix A we consider the methods for a single sine-cosine pair in detail. For convenience, the subscript,  $n$ , is omitted.



## Phase 2

In order to obtain the change in temperature resulting from placement near the surface of a disc maintained at a fixed uniform temperature, a set of 7090 programs has been written to solve the heat equation in three space dimensions with cylindrical symmetry using the standard forward difference approximation to the differential equation. Since the temperature difference is to be obtained, the initial condition in the medium is that the temperature difference is zero. The boundary conditions at the surface, however, must reflect the presence of the radiation interchange with the disc and of the solar radiation (during lunar day) and re-radiation of heat into space. These radiation processes depend on the absolute surface temperature which is calculated by superposition of the temperature difference caused by the disc upon the normal, steady periodic solution found in Phase 1. In Appendix B, the numerical process used is described in detail.

## Results of Phase 1

Using the 7090 program for Phase 1, a number of sets of temperature vs. time profiles were computed at various depths for various choices of physical parameters. Both uniform and layered media were used. The results indicate a definite limitation on the possible values of the parameters for the medium. If the parameters are physically unrealistic, that is, unrealistic based on our knowledge of and assumptions about the lunar surface,

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the solutions quickly result in extremely high or extremely low, even negative, temperatures at depths below the surface. Thus, one can establish, within the limits of our assumptions, the maximum depth to which a surface layer of any chosen composition may extend. Below that maximum depth, another type of layer must be assumed to preserve the stability of the solution. It should be noted that in this problem the solution in an upper layer is independent of the choice of lower layers. Of course, the fact that the solutions are well-behaved in an upper layer does not show that it is possible to select lower layers so that the solution will continue to be well behaved.

The following values for parameters were used in all cases:

1.  $\alpha_v = 0.875$
2.  $q_0 = 1.95 \text{ cal/cm}^2 \text{ min} = 117.0 \text{ cal/cm}^2 \text{ hr.}$
3.  $\Omega = 1.4776 \times 10^{-4} \text{ min}^{-1} = 8.8656 \times 10^{-3} \text{ hr}^{-1}.$
4.  $\epsilon_s = 0.9$
5.  $\sigma = 5.669 \times 10^{-12} \text{ watt/cm}^2 (\text{°K})^4 = 4.878 \times 10^{-9} \text{ cal/cm}^2 \text{ hr } (\text{°K})^4.$

Lambert's Law was assumed to hold for all surfaces.

Three basic types of material were chosen to make up layers. They were:

1. "Dust\*" :  $K = 1.65 \times 10^{-5} \text{ cal/sec cm } \text{°K}$   
 $= 5.94 \times 10^{-2} \text{ cal/hr cm } \text{°K}$   
 $\rho = 2.0 \text{ gram/cm}^3$   
 $c = 0.2 \text{ cal/gram } \text{°K}$

\* Taken to be dust in vacuum.

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2. "Pumice" :  $K = 5.37 \times 10^{-4} \text{ cal/sec cm } ^\circ\text{K}$   
                   $= 1.93 \text{ cal/hr cm } ^\circ\text{K}$   
                   $\rho = 0.6 \text{ gram/cm}^3$   
                   $c = 0.2 \text{ cal/gram } ^\circ\text{K}$
3. "Basalt" :  $K = 5.37 \times 10^{-3} \text{ cal/sec cm } ^\circ\text{K}$   
                   $= 1.93 \times 10^1 \text{ cal/hr cm } ^\circ\text{K}$   
                   $\rho = 2.8 \text{ gram/cm}^3$   
                   $c = 0.2 \text{ cal/gram } ^\circ\text{K}$

Several sets of two-layered cases were run. In addition, a wide range of conductivities was explored using a uniform medium with  $\rho c$  the same as "dust". These may show the maximum depths to which such a top layer may be used.

In order to specifically eliminate certain media from the class of "reasonable" media, improved representation of surface data would be required as indicated below.

### Results of Phase 2

Most of the calculations done in Phase 2 used the radiation boundary condition corresponding to the disc being placed very close to the surface so that edge effects could be ignored. However, if the disc is not assumed to be very close, allowance must be made for loss of energy by radiation from areas under the disc into the sky. This will cause a more rapid perturbation of the uniformity of the radial surface temperature distribution.

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In order to explore this possibility a calculation was made using a radiation condition for a circular disc, parallel to and above the surface. Then, the direct radiation flux from the disc absorbed by the surface is,

$$q_{sh} = \epsilon_s \epsilon_{sh} \sigma T_{sh}^4 \Phi. \quad (5)$$

$\epsilon_s$  is the lunar surface emittance,  $\epsilon_{sh}$  is the disc surface emittance, and  $T_{sh}$  is the temperature of the disc.  $R_{sh}$  is the radius of the disc,  $R_s$  is the radial distance on the surface to the point of interest,  $h$  is the height of the disc above the surface, and,

$$\Phi = \frac{1}{2} \left\{ 1 - \frac{1 + C^2 - B^2}{\sqrt{C^4 + 2C^2(1-B^2) + (1+B^2)^2}} \right\} \quad (6)$$

where  $B = R_{sh}/h$  and  $C = R_s/h^3$ .

For the assumption that the disc is very close (HEIGHT =  $h = 0.0$ ) to the surface, we use,

$$\Phi = 1, R_s < R_{sh}, \text{ and} \quad (7a)$$

$$\Phi = 0, R_s > R_{sh}. \quad (7b)$$

<sup>3</sup>Heat Transfer, Vol. II, Max Jakob, John Wiley and Sons, New York, 1957, p. 11.

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For this case the correct total surface flux for  $R_s < R_{sh}$  is given by

$$q_{net} = E\sigma (T_{sh}^4 - T_{su}^4) \quad (8)$$

where  $T_{su}$  is the surface temperature and

$$E = (1/\epsilon_s + 1/\epsilon_{sh} - 1)^{-1}. \quad (9)$$

Equation (8) includes all reflections and  $E$  is the sum of either of the series,

$$E = \epsilon_s \epsilon_{sh} \sum_{m=0}^{\infty} (1-\epsilon_{sh})^m (1-\epsilon_s)^m \quad (10a)$$

$$= \epsilon_s [1 - (1-\epsilon_{sh}) \epsilon_s \sum_{m=0}^{\infty} (1-\epsilon_{sh})^m (1-\epsilon_s)^m] \quad (10b)$$

where (10a) represents the reflections applied to  $T_{sh}^4$  and (10b) represents those applied to  $T_{su}^4$ .

For  $h \neq 0$ , the reflection considerations are much more complicated. Using the simple form given by (5) would give results closely comparable to  $h = 0$  results if the reflection terms in (10) were neglected. Hence, for  $h = 0$  we used only the leading terms of the series (10), thus,

$$q_{net} = \sigma (\epsilon_s \epsilon_{sh} T_{sh}^4 - \epsilon_s T_{su}^4). \quad (11)$$

Several values of the product  $K_{pc}$  were assigned, and temperature perturbation calculations were made using these values

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for the case  $h = 0$ . These results indicate that it would be possible to experimentally determine the product  $K_{pc}$  for the near-surface lunar material using the apparatus previously described.

The single calculation using  $h \neq 0$  indicates that it would be possible to determine  $K_{pc}$  if the disc were placed at most as far from the surface as 5 cm.

The calculations described above were done for initial time (time of disc placement) = 168.0 hrs. after sunrise. This time is approximately "noon" of the lunar day. In addition, a single calculation was made for "midnight" initial time. All results graphically presented in Appendix D are for "noon" initial time. Although no detailed study of results for "midnight" initial time was made, it is expected that  $K_{pc}$  could also be determined by an experiment at "midnight". These two initial times were the only ones considered since it seemed desirable to choose an initial time at which the true derivative of normal surface temperatures is very small.

#### Comments

The application of the methods described above to an actual experiment would, of course, require additional calculation to improve the accuracy of perturbed temperatures given in Appendix D.

No rigorous analysis of errors has been attempted. In particular, no attempt has been made to estimate the effect of

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using Eqn. (11) rather than Eqn. (8) to express the radiation interchange at the surface.

The surface temperatures used in this work were obtained from a plot entitled, "Temperature Variation of the Lunar Surface", furnished by North American Aviation, Inc. Inaccuracies in  $v(o,t)$  are magnified in the surface flux calculated by Eqn. (2), since  $f(o,t)$  is small compared to either of the terms  $\alpha_v q$  and  $\epsilon_g \sigma [v(o,t)]^4$ . This inaccuracy leads to a physically unrealistic magnification with depth of the higher frequency components of temperature and flux. In the absence of raw temperature data, no attempt was made to estimate the accuracy of the individual Fourier components for temperature or flux. The 20 sine-cosine pairs were necessary to represent the temperature plot used in the calculations. It is expected that Fourier sums calculated from raw data would justify using fewer components which would significantly improve the results of Phase 1, and possibly affect the results of Phase 2.

In an actual experiment, the initial (normal) surface temperature should be measured to allow correction of the astronomical values previously used. Then perturbed temperatures should be measured at a sequence of times up to approximately one-third of an hour. Then an estimate of the product  $K_{pc}$  would be determined from an expanded improved version of Fig. 4 for each perturbed temperature measured. An analysis of the variation of these estimates of  $K_{pc}$  may give some indication of the reliability of the method.

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In general, the results of this work confirm the preliminary design calculations for surface determination of near-surface thermal properties.



APPENDIX ADETERMINATION OF INITIAL CONDITIONS

To determine initial conditions existing at placement of the disc, consider the problem:

$$v_t = x v_{xx} ; 0 < x < \gamma, -\infty < t < \infty, \quad (1a)$$

$$v(0,t) = (2p) \cos \omega t + (2q) \sin \omega t; -\infty < t < \infty \quad (1b)$$

$$f(0,t) = -K \frac{\partial v}{\partial x}(0,t) = (-4K\zeta r) \cos \omega t + (-4K\zeta s) \sin \omega t \quad (1c)$$

$$\gamma, \omega, x, K, p, q, r, s \text{ constant}; \zeta^2 = \frac{1}{2} \omega/x, \omega \neq 0. \quad (1d)$$

$$v(x,t) = [P \cos (\omega t + \zeta x) + Q \sin (\omega t + \zeta x)] e^{\zeta x} \\ + [R \cos (\omega t - \zeta x) + S \sin (\omega t - \zeta x)] e^{-\zeta x}, \quad (2a)$$

satisfies (1a) (Appendix C.I). Hence, if constants P, Q, R, S can be chosen satisfying (1b) and (1c), then  $v(x,t)$  given by (2a) will satisfy (1).

Differentiating (2a) we obtain,

$$f(x,t) = -K \frac{\partial v}{\partial x} = -K\zeta \{ [-P \sin (\omega t + \zeta x) + P \cos (\omega t + \zeta x) \\ + Q \cos (\omega t + \zeta x) + Q \sin (\omega t + \zeta x)] e^{\zeta x} \\ - [-R \sin (\omega t - \zeta x) + R \cos (\omega t - \zeta x) \\ + S \cos (\omega t - \zeta x) + S \sin (\omega t - \zeta x)] e^{-\zeta x} \}. \quad (2b)$$

Requiring (2a) and (2b) to reduce to (1b) and (1c), respectively for  $x = 0$ , we obtain,

$$2p = P + R; 2q = Q + S, \quad (3a)$$

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$$\text{and } 4r = P + Q - R - S; 4s = -P + Q + R - S, \quad (3b)$$

$$\text{or } 2(r+s) = Q - S; 2(r-s) = P - R. \quad (3c)$$

$$\begin{aligned} P &= p + r - s \\ \therefore, Q &= q + r + s \\ R &= p - r + s \\ S &= q - r - s \end{aligned} \quad (4)$$

We wish to write (2) in the form

$$v(x,t) = P_I^*(x) \cos \omega t + Q_I^*(x) \sin \omega t, \quad (5a)$$

$$f(x,t) = -K\zeta \{R_I^*(x) \cos \omega t + S_I^*(x) \sin \omega t\}. \quad (5b)$$

To do so we apply trigonometric addition formulae to (2) obtaining  
(Appendix C.II, C.III)

$$\begin{aligned} P_I^*(x) &= P^*(x) + R^*(x) \\ Q_I^*(x) &= Q^*(x) + S^*(x) \\ R_I^*(x) &= Q^*(x) - R^*(x) + P^*(x) - S^*(x) \\ S_I^*(x) &= Q^*(x) + R^*(x) - P^*(x) - S^*(x), \end{aligned} \quad (6)$$

where

$$\begin{aligned} P^*(x) &= (P \cos \zeta x + Q \sin \zeta x) e^{\zeta x} \\ Q^*(x) &= (-P \sin \zeta x + Q \cos \zeta x) e^{\zeta x} \\ R^*(x) &= (R \cos \zeta x - S \sin \zeta x) e^{-\zeta x} \\ S^*(x) &= (R \sin \zeta x + S \cos \zeta x) e^{-\zeta x} \end{aligned} \quad (7)$$

$$\text{Given data } v(0,t) = P_D \cos \omega t + Q_D \sin \omega t \quad (8a)$$

$$f(0,t) = R_D \cos \omega t + S_D \sin \omega t \quad (8b)$$

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We compute 
$$p = \frac{1}{2} P_D, \quad q = \frac{1}{2} Q_D$$

$$r = -\frac{1}{4K\zeta} R_D, \quad s = -\frac{1}{4K\zeta} S_D$$

(9)

Then,  $P, Q, R, S$  are given by (4). Then, (2) gives temperature and flux throughout  $0 < x < \gamma, -\infty < t < \infty$ . Alternatively, for fixed  $x \in (0, \gamma)$ , (5) gives temperature and flux for all time.

In particular, if  $x = \gamma$  represents an interface between two layers;

$$P_I^*(\gamma), Q_I^*(\gamma), -K_1 \zeta_1 R_I^*(\gamma), -K_1 \zeta_1 S_I^*(\gamma),$$

(10)

respectively play the role of

$$P_D, Q_D, R_D, S_D$$

for computation in a second layer by reduction to the previous case where  $K_1$  is the conductivity appropriate to the first layer.

Distance,  $x$ , in the second layer will be measured from the interface rather than from the top of the first layer. Clearly, this process may be continued for as many layers as desired.

The process has two parts, viz.

- 1° We cross an interface, i.e., we apply (4) to data representing input at the top of a layer.
- 2° We use coefficients so obtained to carry ourselves into or through a layer by application (6) and (7) to obtain (5).

In the computer, (7a) for example, may be difficult to

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evaluate for the desired  $x$  due to overflow on  $e^{\xi x}$ ; even when  $P^*(x)$  itself is not over-large. Hence, it may be desirable to pretend the existence of certain interfaces, i.e., we may represent a single thick layer by several thinner layers of identical physical properties to guarantee smallness of  $\xi x$  throughout every layer (recall  $x$  is always measured from the top of the layer of interest).

In the special case,  $\omega = 0$ , the above formulation does not apply. That is, given data

$$v(0,t) = a; \quad f(0,t) = -Kb \quad (1b1)$$

$$\text{we find the solution} \quad v(x,t) = a + bx \quad (1c1)$$

$$f(x,t) = -Kb \quad (2a1)$$

$$(2b1)$$

Now, if we are given data

$$v(0,t) = \sum_{n=0}^N \{P_{D_n} \cos \omega_n t + Q_{D_n} \sin \omega_n t\} \quad (8a1)$$

$$f(0,t) = \sum_{n=0}^N \{R_{D_n} \cos \omega_n t + S_{D_n} \sin \omega_n t\}, \quad (8a2)$$

$$\omega_n = n\Omega$$

We simply apply the above results by superposition to obtain our solution. Note that  $a = P_{D_0}$ ,  $-Kb = R_{D_0}$ .

APPENDIX BDETERMINATION OF TEMPERATURE PERTURBATIONS

In Appendix A we dealt with the determination of normal lunar temperatures, i.e., those temperatures which exist in the absence of disturbance by man-made equipment.

These normal temperatures satisfy the heat equation in one space dimension. Since this temperature is known for all time and independent of radial distance,  $r$ , measured from the center of the disc and azimuth,  $\theta$ , the following equations are satisfied:

$$\frac{\partial v}{\partial t} = \kappa \left\{ \frac{\partial^2 v}{\partial x^2} + \frac{1}{r} \frac{\partial v}{\partial r} + \frac{\partial^2 v}{\partial r^2} \right\}; \quad r, x > 0, \quad -\infty < t < \infty \quad (1a)$$

$$\frac{\partial v}{\partial r}(x, r, t) \equiv 0, \text{ in particular } \frac{\partial v}{\partial r}(x, 0, t) = 0; \quad t, x > 0 \quad (1b)$$

$$\frac{\partial v}{\partial x}(0, r, t) - \text{known for all } t. \quad (1c)$$

$$v(x, r, 0) = v(x, r, 0) - \text{known by Phase 1} \quad (1d)$$

$$v(0, r, t) - \text{known from all } t \quad (1e)$$

where time,  $t = 0$ , is naturally taken as the time of placement of the disc.

Now after placement of the disc, i.e., for positive  $t$ , the above relations will not describe the temperature. Rather, the temperature  $\mu(x, r, t)$  will satisfy:

$$\frac{\partial \mu}{\partial t} = \kappa \left\{ \frac{\partial^2 \mu}{\partial x^2} + \frac{1}{r} \frac{\partial \mu}{\partial r} + \frac{\partial^2 \mu}{\partial r^2} \right\}; \quad x, r, t > 0, \quad (1a1)$$

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$$\frac{\partial \mu}{\partial r}(x, 0, t) = 0; \quad t, x > 0 \text{ by cylindrical symmetry,} \quad (1b1)$$

$$-K \frac{\partial \mu}{\partial x}(0, r, t) = F(r, t) - \eta[\mu(0, r, t)]^4; \quad r, t > 0, \quad (1c1)$$

$$\mu(x, r, 0) = v(x, r, 0). \quad (1d1)$$

Since  $v(x, r, t)$  is known in closed form from Phase 1, we consider for simplicity in calculation:

$$w(x, r, t) = \mu(x, r, t) - v(x, r, t). \quad (2)$$

From the above definition of  $\mu$  and  $v$ , we have

$$\frac{\partial w}{\partial t} = x \left\{ \frac{\partial^2 w}{\partial x^2} + \frac{1}{r} \frac{\partial w}{\partial r} + \frac{\partial^2 w}{\partial r^2} \right\}; \quad x, r, t > 0 \quad (1a2)$$

$$\frac{\partial w}{\partial r}(x, 0, t) = 0; \quad t, x > 0 \quad (1b2)$$

$$-K \frac{\partial w}{\partial x}(0, r, t) = F(r, t) - \eta \mu^4(0, r, t) + K \frac{\partial v}{\partial x}(0, r, t); \quad r, t > 0, \quad (1c2)$$

$$w(x, r, 0) \equiv 0. \quad (1d2)$$

Then,  $w(x, r, t)$  represents perturbation of temperature from normal due to the placing of the disc.  $w(x, r, t)$  was computed numerically as described below:

Let  $Z_{ijn} = Z(i\Delta x, j\Delta r, n\Delta t)$  represent the approximation to  $w(i\Delta x, j\Delta r, n\Delta t)$  obtained below. Let the operator  $L$  be defined by

$$LZ_{ijn} = (\Delta x)^2 \left\{ \Delta_x^2 Z_{ijn} + \frac{Z_{i,j+1,n} - Z_{i,j-1,n}}{2j(\Delta r)^2} + \Delta_r^2 Z_{ijn} \right\} \quad (3)$$

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where

$$\Delta_x^2 Z_{ijn} = (Z_{i+1,j,n} - 2Z_{i,j,n} + Z_{i-1,j,n})/(\Delta x)^2. \quad (4)$$

It seems desirable to take  $\Delta x = \Delta r$  and we assume this to be done.

We define  $Z_{ijn}$  by the difference equations (Appendix C.IV)

$$Z_{i,j,n+1} = Z_{ijn} + \frac{x\Delta t}{(\Delta x)^2} LZ_{ijn}; \quad n \geq 0; \quad x, r > 0, \quad (1a3)$$

$$Z_{i,o,n} = Z_{i,1,n}; \quad n \geq 0, \quad i > 0, \quad (1b3)$$

$$Z_{o,j,n} = \phi_{o,j,n} - v_{o,j,n}; \quad n \geq 0, \quad j > 0 \quad (1c3)$$

where  $\phi_{o,j,n}$  is the positive root of

$$\left(\frac{2}{3} \frac{\eta}{K} \Delta x\right) \phi^4 + \phi - \left\{v_{o,j,n} + \frac{1}{3} [4Z_{1,j,n} - Z_{2jn} + 2\Delta x \left(-\frac{F_{j,n}}{K} + \frac{\partial v}{\partial x} o_{j,n}\right)]\right\} = 0 \quad (5)$$

$$Z_{1,j,o} = 0. \quad (1d3)$$

Given values at  $t = n\Delta t$ , we compute

$Z_{1,j,n+1}$  by (1a3);  $ij = 1, 2, \dots$

$Z_{i,o,n+1}$  by (1b3);  $i = 1, 2, \dots$

$Z_{o,j,n+1}$  by (1c3);  $j = 1, 2, \dots$

$Z_{o,o,n+1}$  could be computed by either (1b3) or (1c3), but is immaterial since it is not used for progress to the next time level. However, it might be of interest to compare these two values of  $Z_{o,o,n}$  for large  $n$ .

-B4-

Algebraic Problem: Denote  $A = \frac{2\eta}{3K} \Delta x$  and  $B_{jn} = v_{ojn} + \frac{1}{3} [4Z_{1,j,n} - Z_{2jn} + 2\Delta x \frac{F_{jn}}{K} + \frac{\partial v}{\partial x} o_{jn}]$ . (6)

Now, (5) has precisely one positive root if  $B_{jn} > 0$  and no positive root if  $B_{jn} \leq 0$ , since the function  $g(\phi) = A\phi^4 + \phi - B$  is negative at  $\phi = 0$  for positive  $B$  and positive for negative  $B$ ;  $g'(\phi) = 4A\phi^3 + 1$  is positive for all positive  $\phi$ ; and  $g(\phi)$  is positive for sufficiently large  $\phi$ .

Thus, we have  $Z_{ijn}$  completely defined. For stable choices of  $\Delta x$ ,  $\Delta r$ ,  $\Delta t$ ,  $B$  is always positive. In fact,  $v_{o,j,n}$  is normally the dominant term in  $B$  for sufficiently small  $\Delta x$ .

Although  $\mu_{ijn}$ , especially  $\mu_{ojn}$ , is desirable output,  $Z_{ijn}$  has obvious computer storage advantages. In fact,  $Z_{ijn}$  at time  $t = n\Delta t$  is defined to be zero for all sufficiently large  $i$  and  $j$ . Thus,  $Z_{ijn}$  requires much less storage than  $\mu_{ijn}$ . Even so, the magnitude of the problem which may be solved without use of magnetic tape for interim storage is rather small. We made no provision for this interim storage; hence, our results will be given for times of, at most, one hour.

Throughout the computations we require  $\Delta x = \Delta r$  and  $\Delta t$  to be chosen to satisfy

$$\frac{x\Delta t}{(\Delta x)^2} < \frac{1}{4} \quad (7)$$

to maintain numerical stability<sup>1</sup>.

<sup>1</sup>Finite Difference Methods for Partial Differential Equations, G. E. Forsythe, W. R. Wasow, Wiley, 1960



-B5-

Most references give  $1/2$  where we have given  $1/4$ . However, the proper expression depends on the number of space dimensions actually appearing in the heat equation.  $1/2$  is correct for one-space dimension, while  $1/4$  is correct for two dimensions ( $x$  and  $r$ ).

The above considerations are altered by the presence of an interface between layers. If  $\Delta x$  is a depth corresponding to an interface we determine  $Z_{i,j,n+1}$  by the expression

$$(K_1 + K_2) Z_{i,j,n+1} = K_2 Z_{i+1,j,n+1} + K_1 Z_{i-1,j,n+1}, \quad (8)$$

which approximates

$$K_1 \frac{\partial w}{\partial x}(x=0, r, t) = K_2 \frac{\partial w}{\partial x}(x=0, r, t) \quad (9a)$$

$$w(x=0, r, t) = w(x=0, r, t), \quad (9b)$$

where  $K_1$  and  $K_2$  are conductivities of the first and second layers, respectively. Otherwise, relations (1a3), (1b3), (1c3), and (1d3) apply. In this case the diffusivity,  $\alpha$ , appearing in the stability requirement (7) is the largest diffusivity encountered in the problem.

APPENDIX CMATHEMATICAL NOTES

I.  $v(x,t)$  given by (A2a) satisfies (A1a):

$$\begin{aligned} x \frac{\partial^2}{\partial x^2} \{e^{\pm \zeta x} \cos (\omega t \pm \zeta x)\} &= x(\pm \zeta) \frac{\partial}{\partial x} \{e^{\pm \zeta x} [-\sin (\omega t \pm \zeta x) \\ &+ \cos (\omega t \pm \zeta x)]\} = x \zeta^2 e^{\pm \zeta x} \{-\sin (\omega t \pm \zeta x) + \cos (\omega t \pm \zeta x) \\ &- \cos (\omega t \pm \zeta x) - \sin (\omega t \pm \zeta x)\} = \frac{1}{2} \omega e^{\pm \zeta x} \{-2 \sin (\omega t \pm \zeta x)\} \\ &= -\omega e^{\pm \zeta x} \sin (\omega t \pm \zeta x) = \frac{\partial}{\partial t} \{e^{\pm \zeta x} \cos (\omega t \pm \zeta x)\}. \end{aligned}$$

Since  $\sin (\omega t \pm \zeta x) = \cos (\omega t \pm \zeta x - \pi/2)$ ,  $e^{\pm \zeta x} \sin (\omega t \pm \zeta x)$  also satisfies (A1a); hence,  $v(x,t)$  satisfies (A1a).

II. Addition formulae applied to  $v(x,t)$ :

$$\begin{aligned} \text{By (A2a), } v(x,t) &= \{[P(\cos \omega t \cos \zeta x - \sin \omega t \sin \zeta x) + Q(\sin \omega t \\ &\cos \zeta x + \cos \omega t \sin \zeta x)]e^{\zeta x} + [R(\cos \omega t \\ &\cos \zeta x + \sin \omega t \sin \zeta x) + S(\sin \omega t \cos \zeta x \\ &- \cos \omega t \sin \zeta x)]e^{-\zeta x}\} \\ &= \{[(P \cos \zeta x + Q \sin \zeta x)e^{\zeta x} + (R \cos \zeta x \\ &- S \sin \zeta x)e^{-\zeta x}] \cos \omega t + [(-P \sin \zeta x \\ &+ Q \cos \zeta x)e^{\zeta x} + (R \sin \zeta x + S \cos \zeta x)e^{-\zeta x}] \\ &\sin \omega t\} \\ &= [P^*(x) + R^*(x)] \cos \omega t + [Q^*(x) + S^*(x)] \\ &\sin \omega t. \end{aligned}$$

-C2-

III. Addition Formulae applied to  $f(x,t)$ :

We obtain the desired result by differentiating  $v(x,t)$  since application of addition formulae is tedious. From (A5a) and (A7),

$$f(x,t) = -K \frac{\partial v}{\partial x} = -K \zeta \{ [Q^*(x) - R^*(x) + P^*(x) - S^*(x)] \cos \omega t \\ + [Q^*(x) + R^*(x) - P^*(x) - S^*(x)] \sin \omega t \}$$

since

$$\frac{dP^*}{dx} = \zeta [P^*(x) + Q^*(x)],$$

$$\frac{dQ^*}{dx} = \zeta [Q^*(x) - P^*(x)],$$

$$\frac{dR^*}{dx} = \zeta [-R^*(x) - S^*(x)], \text{ and}$$

$$\frac{dS^*}{dx} = \zeta [-S^*(x) + R^*(x)].$$

IV. Divided Differences: Let  $w = w(x,r,t)$  and  $w_{ijn} = w(i\Delta x, j\Delta r, n\Delta t)$ ; then

$$w_{1,j+1,n} = w_{ijn} + \Delta r \frac{\partial w}{\partial r} i j n + \frac{(\Delta r)^2}{2} \frac{\partial^2 w}{\partial r^2} i j n + \frac{(\Delta r)^3}{6} \frac{\partial^3 w}{\partial r^3} i j n \\ + O((\Delta r)^4),$$

$$w_{1,j-1,n} = w_{ijn} - \Delta r \frac{\partial w}{\partial r} i j n + \frac{(\Delta r)^2}{2} \frac{\partial^2 w}{\partial r^2} i j n - \frac{(\Delta r)^3}{6} \frac{\partial^3 w}{\partial r^3} i j n \\ + O((\Delta r)^4).$$

$$\therefore w_{1,j+1,n} - w_{1,j-1,n} = 2\Delta r \frac{\partial w}{\partial r} i j n + O((\Delta r)^3).$$

$$\text{Also, } w_{1,j+1,n} + w_{1,j-1,n} = 2w_{ijn} + (\Delta r)^2 \frac{\partial^2 w}{\partial r^2} i j n + O((\Delta r)^4).$$

-C3-

Similarly,  $w_{i+1,j,n} + w_{i-1,j,n} = 2w_{ijn} + (\Delta x)^2 \frac{\partial^2 w}{\partial x^2} i_{jn} + O((\Delta x)^4)$ .

$$w_{2jn} = w_{ojn} + 2\Delta x \frac{\partial w}{\partial x} o_{jn} + \frac{4(\Delta x)^2}{2} \frac{\partial^2 w}{\partial x^2} o_{jn} + O((\Delta x)^3),$$

$$w_{1jn} = w_{ojn} + \Delta x \frac{\partial w}{\partial x} o_{jn} + \frac{(\Delta x)^2}{2} \frac{\partial^2 w}{\partial x^2} o_{jn} + O((\Delta x)^3).$$

$$\therefore w_{2jn} - 4w_{1jn} = -3w_{ojn} - 2\Delta x \frac{\partial w}{\partial x} o_{jn} + O((\Delta x)^3).$$

$$\begin{aligned} Lw_{ijn} &= (\Delta x)^2 \left\{ \Delta_x^2 w_{ijn} + \frac{w_{i,j+1,n} - w_{i,j-1,n}}{2r \Delta r} + \Delta_r^2 w_{ijn} \right\} \\ &= (\Delta x)^2 \left\{ \frac{\partial^2 w}{\partial x^2} i_{jn} + O((\Delta x)^2) + \frac{1}{r} \left[ \frac{\partial w}{\partial r} i_{jn} + O((\Delta r)^2) \right] \right. \\ &\quad \left. + \frac{\partial^2 w}{\partial r^2} i_{jn} + O((\Delta r)^2) \right\} \end{aligned}$$

$$= (\Delta x)^2 \{ w_{xx} + r^{-1} w_r + w_{rr} \} + O((\Delta x)^4) + \frac{(\Delta r)^2 (\Delta x)^2}{r} + (\Delta x)^2 (\Delta r)^2$$

$$= \frac{(\Delta x)^2}{r} w_t + O((\Delta x)^4) + \frac{(\Delta r)^2 (\Delta x)^2}{r} + (\Delta x)^2 (\Delta r)^2$$

$$= \frac{(\Delta x)^2}{r} \left\{ \frac{w_{i,j,n+1} - w_{i,j,n}}{\Delta t} + O(\Delta t) \right\} + O((\Delta x)^4) + \frac{(\Delta r)^2 (\Delta x)^2}{r} + (\Delta x)^2 (\Delta r)^2.$$

$$\therefore \frac{\kappa \Delta t}{(\Delta x)^2} Lw_{ijn} = w_{i,j,n+1} - w_{ijn} + O((\Delta t)^2) + \Delta t (\Delta x)^2 + \frac{\Delta t (\Delta r)^2}{r} + \Delta t (\Delta r)^2.$$

APPENDIX DSUMMARY OF NUMERICAL RESULTS

All calculations were based on the temperature and flux Fourier coefficients given below. The flux coefficients give flux units cal/sec cm<sup>2</sup>, hence, the 7090 programming multiples flux coefficients by 3600.0 to obtain flux units cal/hr cm<sup>2</sup>. Since flux depends on the parameters shown above (Eqn. 4), it would be necessary to recompute flux coefficients if the parameters were to be altered.

Time profiles of normal temperatures in "dust" are given below for several depths (Fig. 1). For other media, the general form of the time profiles is very similar. This is indicated by application of the transformation

$$v(x,t) \equiv V(\xi,t) \text{ where } \xi = \int_0^x \frac{d\alpha}{K(\alpha)} \quad (1)$$

to Equations A1, obtaining

$$\frac{\partial V}{\partial t} = \frac{x}{K^2} \frac{\partial^2 V}{\partial \xi^2} = \frac{1}{K\rho c} \frac{\partial^2 V}{\partial \xi^2} \quad (2a)$$

$$V(0,t) = v(0,t) \quad (2b)$$

$$f(0,t) = - \frac{\partial V}{\partial \xi} (0,t). \quad (2c)$$

Thus, the time profiles in the  $(\xi,t)$  coordinate system depend only on the product  $K\rho c$ .

-2-

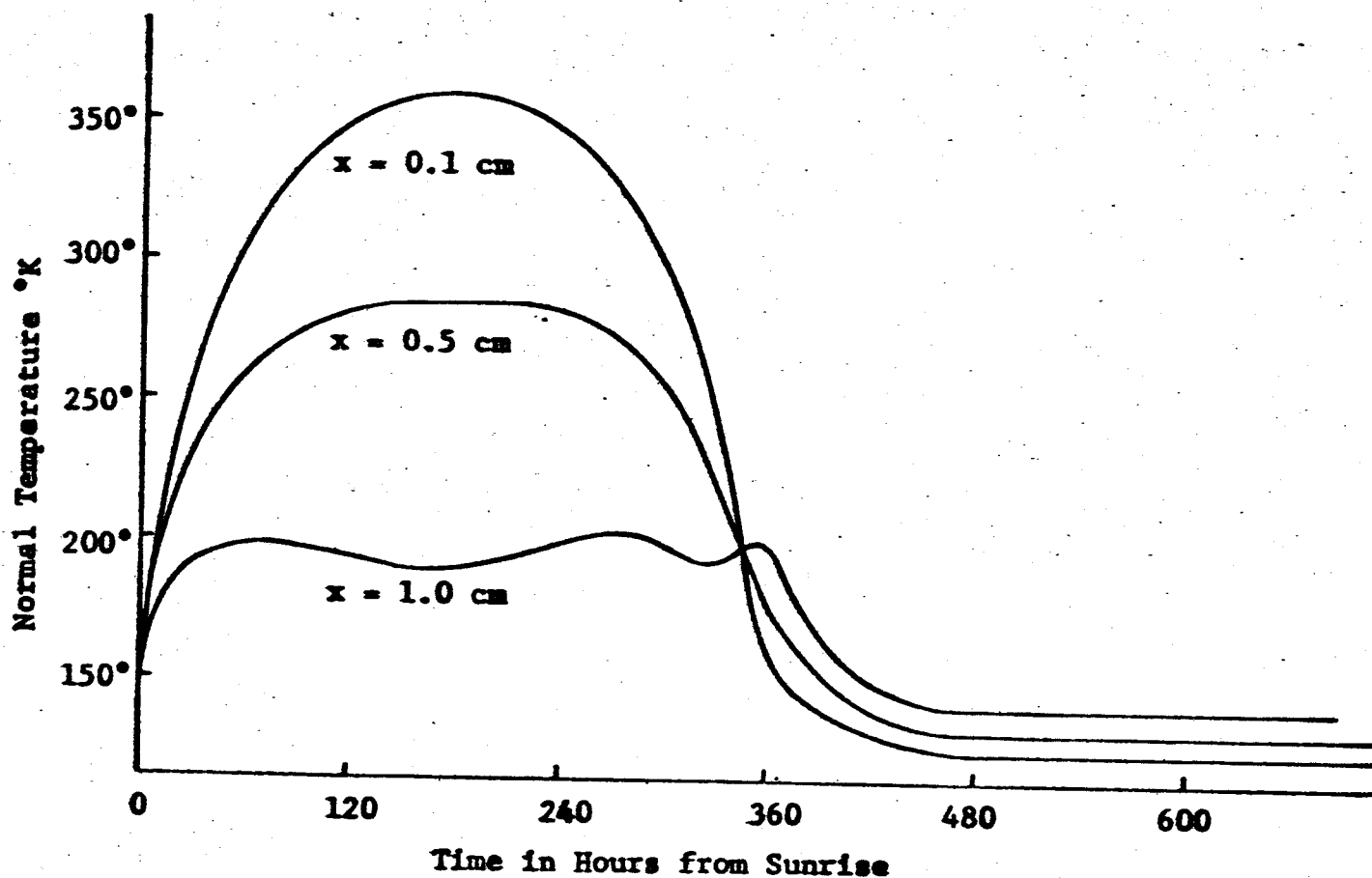
Curves representing perturbation of temperatures for several values of  $K_{pc}$  are presented in Fig. 2. Since these curves were obtained using height of disc = HEIGHT = 0.0, no angular effect (radiation into space) is considered in these curves. For comparison, one case was computed for HEIGHT = 5.0 cm. The results are compared in Fig. 3. Note that the results indicate that the radial variation of perturbed temperature for a specific time is small over a radius of 4 or 5 cm. Thus, the perturbed temperature is nearly independent of radius over the area of view of the radiometer.

For sufficiently small time so that the radius of the disc has little effect,  $(K_{pc})^{-1/2}$  can be considered a function of the temperature perturbations. Thus, if the surface temperature perturbation,  $w(o,o,t)$ , is known for a definite value of  $t$ ,  $(K_{pc})^{-1/2}$  can be estimated as indicated in Fig. 4.

Since the Phase 2 calculations used finite-difference techniques, there is some error in the temperature perturbations obtained. Although no rigorous analysis of errors has been made, an indication of errors is obtained by comparison of results for successively smaller  $\Delta t$  and  $\Delta r$ . The temperature perturbations obtained in "dust" using  $\Delta x = \Delta r = 0.1$  cm, 0.05 cm, 0.025 cm, successively, differ by 1°K or less where  $\Delta t$  is chosen to maintain numerical stability.

TABLE I  
FOURIER COEFFICIENTS

N	Temp. Sine	Temp. Cos	Flux Sine	Flux Cos	
DATA					
00	.00000000E	00 .22626791E	03 .00000000E	00 .91183791E-03	LUNARI
01	.14474085E	03-.52482391E	01 .18406930E-02	.11056784E-03	LUNARI
02	.16495898E	01-.34031582E	02 .20146566E-04	.63172640E-03	LUNARI
03	.22540356E	02-.33538482E	01 .21957966E-03	.44425867E-04	LUNARI
04	.20174133E	01-.13186581E	02-.13201978E-04	.21278920E-03	LUNARI
05	.84914122E	01-.19856164E	01 .11512204E-03	.36994564E-04	LUNARI
06	.15525196E	01-.76848104E	01-.90358427E-05	.13197542E-03	LUNARI
07	.43425736E	01-.12548302E	01 .59816251E-04	.28458076E-04	LUNARI
08	.11317164E	01-.52154533E	01-.26333029E-05	.85040427E-04	LUNARI
09	.26702678E	01-.90534500E	00 .37523709E-04	.21839815E-04	LUNARI
10	.78925178E	00-.36548020E	01 .31907033E-05	.62879271E-04	LUNARI
11	.19471356E	01-.79391820E	00 .23770127E-04	.16761790E-04	LUNARI
12	.62367758E	00-.31022627E	01 .67389333E-05	.45032700E-04	LUNARI
13	.13407852E	01-.79024720E	00 .19345902E-04	.11874344E-04	LUNARI
14	.37446378E	00-.26508531E	01 .94848804E-05	.37742956E-04	LUNARI
15	.90879622E	00-.69753089E	00 .13012129E-04	.51275089E-05	LUNARI
16	.22515756E	00-.21995335E	01 .62070178E-05	.31778413E-04	LUNARI
17	.50279911E	00-.62417482E	00 .90380011E-05	.48891571E-05	LUNARI
18	.51452000E-01	.17863406E	01 .74823758E-05	.25173496E-04	LUNARI
19	.34896822E	00-.47163624E	00 .90437187E-05	.17235374E-05	LUNARI
20	.23009778E-01	.14265451E	01 .60155256E-05	.24689393E-04	LUNARI
-01	.00000000E	00 .00000000E	00 .00000000E	00 .00000000E	LUNARI

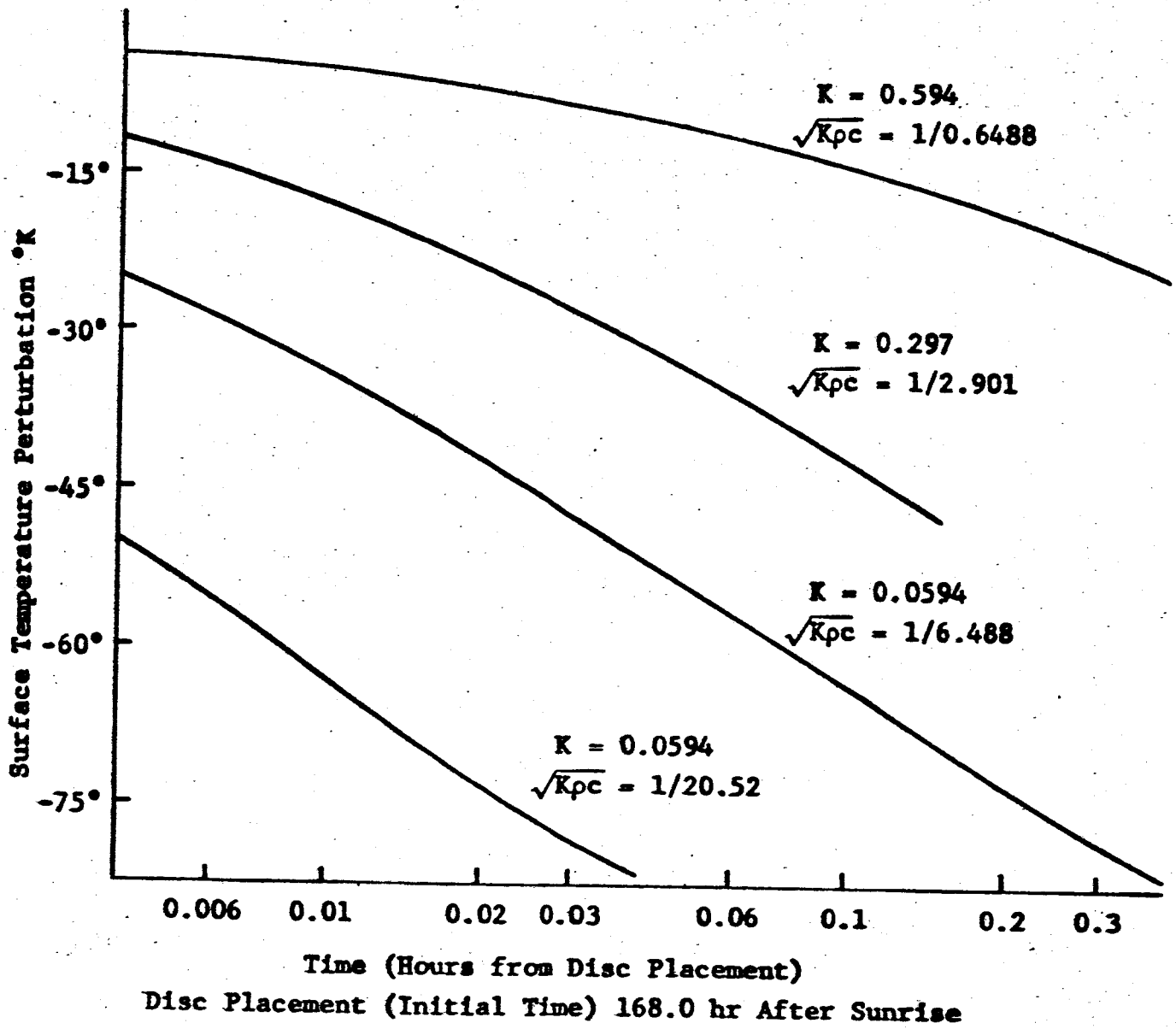


TIME PROFILES OF NORMAL TEMPERATURES IN "DUST"

FIGURE 1

1:794.38-D4  
1:794-86

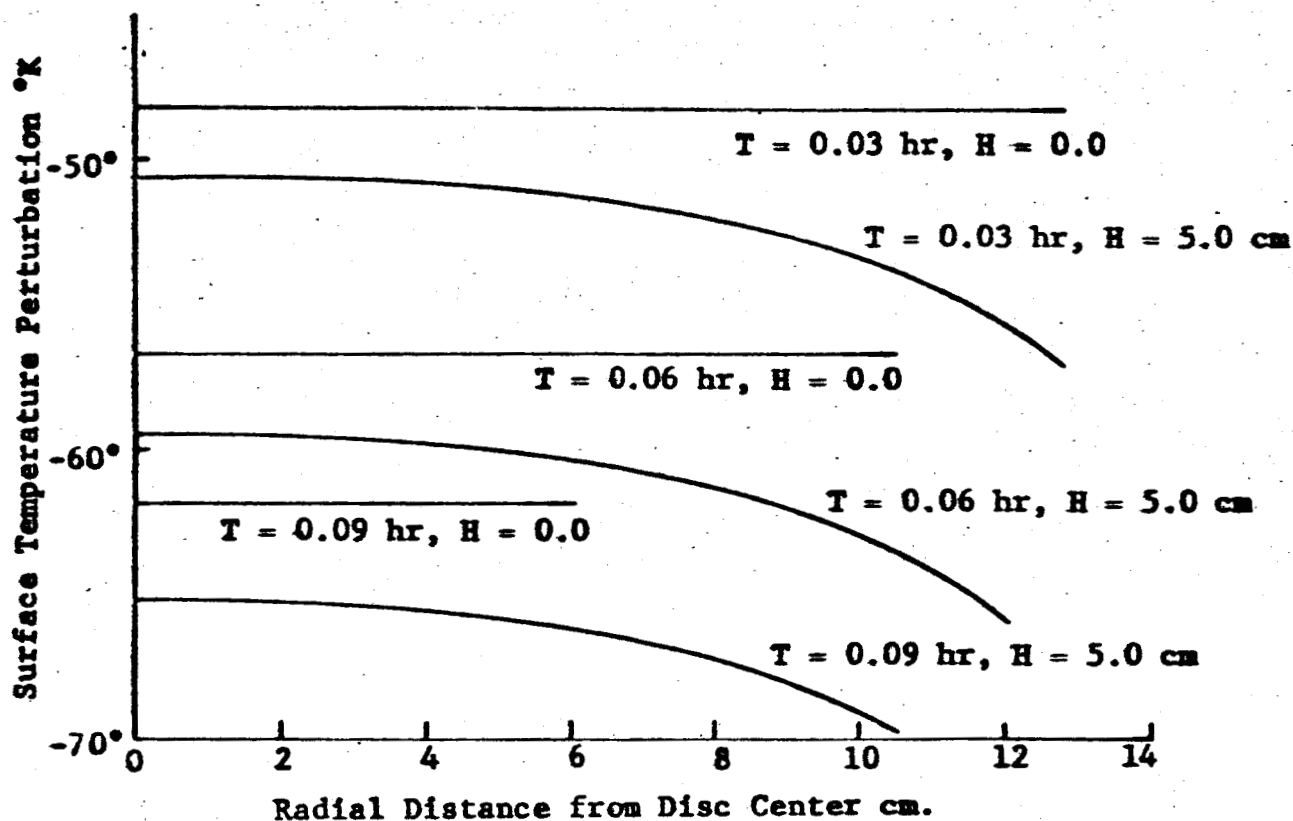




SURFACE TEMPERATURE PERTURBATION FOR VARIOUS UNIFORM MEDIA

FIGURE 2

1:794.38-D5  
1:794-87



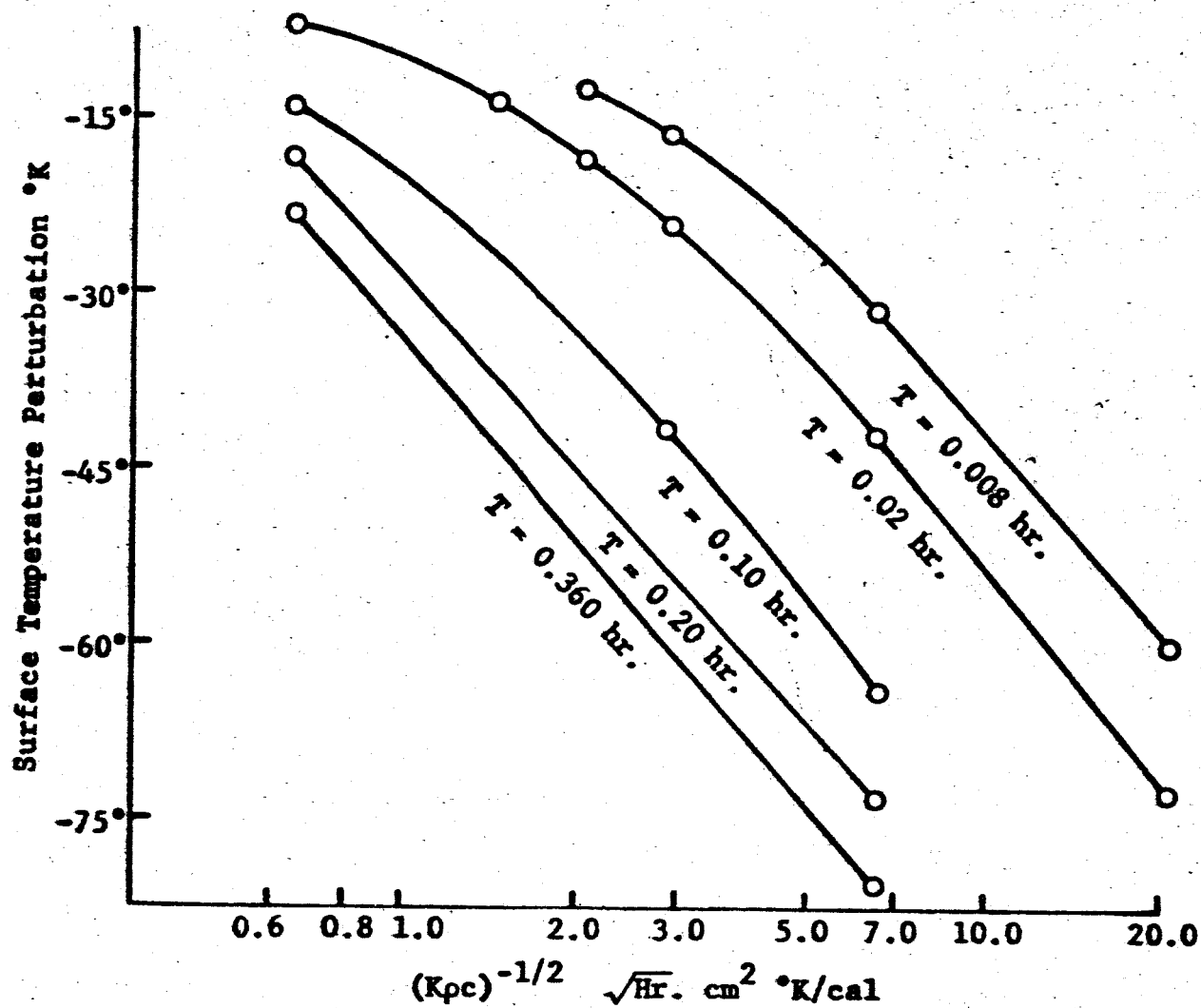
Disc Placement (Initial Time) 168.0 hr. after Sunrise  
Results for "Dust" using  $\Delta x = \Delta r = 0.1 \text{ cm}$ .

Radius of Disc 15.24 cm.

EFFECT OF HEIGHT (H) VARIATION ON PERTURBATIONS VS.  
RADIAL DISTANCE

FIGURE 3

1:794.38-D6  
1:794-88



Disc Placement (Initial Time) 168.0 hr.  
After Sunrise

Results using  $\Delta x = \Delta r = 0.05 \text{ cm}$ . Radius of Disc = 15.24 cm.  
Height of Disc = 0.0

DETERMINATION OF  $\sqrt{K\rho c}$  FROM PERTURBED TEMPERATURES

FIGURE 4

APPENDIX E7090 PROGRAMS

All programming and computation for the preparation of this report was done on Texaco IBM 7090, Houston, Texas with the exception of the determination of Fourier coefficients, which was accomplished on the Texaco Elecom 120A, Bellaire, Texas.

Fortran source listings of this 7090 programming appear below.

These programs operated effectively in the solution of those problems undertaken in the preparation of this report. However, no claim to programming perfection is made. In fact, certain improvements are known to be possible.

The Fortran statement, COMMON, does not appear in Phase 1. This statement is used extensively in Phase 2. In fact, most portions of Phase 2 contain a complete set of Phase 2 COMMON and DIMENSION statements regardless of necessity. These statements are designed to be useful in Phase 1 if Phase 1 were appropriately revised.

Storage requirements listed below are given as decimal numbers. A glossary of variable names is given on Page 40 of this Appendix.

-2-

**Entry Name** - LUNAR1

**Category** - Main Program - Phase 1

**Purpose** - To read data and guide output of normal temperatures as specified by input control.

**Arguments** - None (Main Program)

**Unusual Cautions** - None

**Description** - After input has been read, calculation is made as described in Appendix A using appropriate subroutines as needed.

**Lower Memory Requirements** - 1613

**Transfer Vector** - (FPT), (STH), (FIL), (TSH), (RTN),  
FILET, FILEX, EXIT

**Common Requirements** - None

```

*      LIST
*      LABEL
CLUNARI ALPHABET FOR LABEL
90  FORMAT (1X,1I3,4E14.8)
92  FORMAT (1X,1I3,2E14.8,1I4)
93  FORMAT (1H1,27HLUNARI TEMPERATURE PROFILES)
94  FORMAT (1H0,18HPHYSICAL CONSTANTS)
95  FORMAT (1H,1I3,4E14.8)
96  FORMAT (1H0,17HTIME PROFILE. X=,E14.8,12H DELTA T =,E14.8)
97  FORMAT (1H0,18HDEPTH PROFILE. T=,E14.8,12H DELTA X =,E14.8)
    DIMENSION N(25),CP(25),CQ(25),CR(25),CS(25),DUSE(5),L(5)
    DIMENSION ANS(999),XP(25),XQ(25),XR(25),XS(25),CON(5),THIC(5)
    DIMENSION DP(25),DQ(25),DR(25),DS(25),RHO(5),SPEH(5)
    K=0
    WRITE OUTPUT TAPE 3,93
1   K=K+1
    READ INPUT TAPE 2, 90, N(K), DQ(K), DP(K), DS(K), DR(K)
    DS(K)=3600.0*DS(K)
    DR(K)=3600.0*DR(K)
    IF (N(K)) 2,1,1
2   K=0
    WRITE OUTPUT TAPE 3,94
3   K=K+1
    READ INPUT TAPE 2,90,L(K),CON(K),THIC(K),RHO(K),SPEH(K)
    WRITE OUTPUT TAPE 3,95,L(K),CON(K),THIC(K),RHO(K),SPEH(K)
    IF (L(K)) 8,9,11
9   DUSE(K)=CON(K)/(RHO(K)*SPEH(K))
    GO TO 3
11  DUSE(K)=CON(K)/(RHO(K)*SPEH(K))
4   READ INPUT TAPE 2,92,KFIL, FIX, SPACE, NREP
    DO 10 I=1, 25, 1
    CP(I) = DP(I)
    CQ(I) = DQ(I)
    CR(I) = DR(I)
    CS(I) = DS(I)
10  CONTINUE
    IF (KFIL) 5,6,7
5   GO TO 2
6   WRITE OUTPUT TAPE 3,96, FIX, SPACE
    CALL FILET (CP,CQ,CR,CS,XP,XQ,XR,XS,DUSE,CON,THIC,
1      ANS,N,L, FIX, SPACE, NREP)
    GO TO 4
7   WRITE OUTPUT TAPE 3,97, FIX, SPACE
    CALL FILEX (CP,CQ,CR,CS,XP,XQ,XR,XS,DUSE,CON,THIC,
1      ANS,N,L, FIX, SPACE, NREP)
    GO TO 4
8   CALL EXIT
    END

```

-E4-

Entry Name	- FILET
Category	- Subroutine - Phase 1 only.
Purpose	- To calculate and report as output a profile of normal temperatures vs. time for fixed depth.
Arguments	<p>- CP, CQ, CR, CS, XP, XQ, XR, XS, DUSE, CON, THIC, ANS, N, L, FIX, SPACE, NREP</p> <p>L : The vector L(M) which indicates type of definition for thickness of each layer M,</p> <p>FIX : The fixed depth for the profile,</p> <p>SPACE: Time between successive values in the profile,</p> <p>NREP : Number of values in the profile beginning with value at T = 0.</p> <p>All other arguments appear in the glossary.</p>
Error Exits	- #5 - 7090 Failure
Description	- Using the subroutines CROSS and CARRY, FILET computes a time profile as described in Appendix A.
Lower Memory Requirements	- 372
Transfer Vector	- ERRORQ, CROSS, CARRY, SIN, COS, (STH), (FIL).
Common Requirements	- None

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```

* LIST
* LABEL
SUBROUTINE FILET (CP,CQ,CR,CS,XP,XQ,XR,XS,DUSE,CON,
1 THIC,ANS,N,L,FIX,SPACE,NREP)
90 FORMAT (1H,10E12.5)
DIMENSION N(25),CP(25),CQ(25),CR(25),CS(25),DUSE(5),L(5)
DIMENSION ANS(999),XP(25),XQ(25),XR(25),XS(25),CON(5),THIC(5)
NREP=NREP
DO 1 I=1,NREP,1
ANS(I)=0.0
1 CONTINUE
K=0
2 M=0
K=K+1
IF (N(K)) 9,8,8
8 X=FIX
TERM=N(K)
3 M=M+1
XCAR=X
IF (L(M)) 4,5,6
4 CALL ERRORQ (5,2)
5 IF (X-THIC(M)) 6,7,7
7 XCAR=THIC(M)
6 CALL CROSS (XP(K),XQ(K),XR(K),XS(K),CP(K),CQ(K),CR(K),CS(K),
1 DUSE(M),CON(M),EGA,ZETA,TERM)
CALL CARRY (XP(K),XQ(K),XR(K),XS(K),CP(K),CQ(K),CR(K),CS(K),
1 CON(M),EGA,ZETA,XCAR)
IF (L(M)) 4,13,11
13 IF (X-THIC(M)) 11,11,10
10 X=X-THIC(M)
GO TO 3
11 TAU=0.0
DO 12 I=1,NREP,1
ANS(I)=ANS(I)+CP(K)*COSF(EGA*TAU)+CQ(K)*SINF(EGA*TAU)
TAU=TAU+SPACE
12 CONTINUE
GO TO 2
9 DO 14 I=1,NREP,10
K=I+9
WRITE OUTPUT TAPE 3,90,(ANS(J),J=I,K)
14 CONTINUE
RETURN
END

```



-E6-

Entry Name	- FILET (modified for plot of output)
Purpose	- Same as FILET (E4).
Category	- Subroutine - Phase 1 only.
Arguments	<p>- CP, CQ, CR, CS, XP, XQ, XR, XS, DUSE, CON, THIC, ANS, N, L. FIX, SPACE, NREP</p> <p>L : The vector L(M) which indicates type of definition for thickness of each layer M,</p> <p>FIX : The fixed depth for the profile,</p> <p>SPACE: Time between successive values in the profile,</p> <p>NREP : Number of values in the profile beginning with value at T = 0.</p> <p>All other arguments appear in the glossary.</p>
Error Exits	- #5 - 7090 failure.
Description	- Using the subroutines CROSS and CARRY, FILET computes a time profile as described in Appendix A.
Name of Deck	- FILTPL - the output is plotted by UMPLOT, University of Michigan Plot, available through SHARE.
Lower Memory Requirements	- 2300 (approximately)
Transfer Vector	- ERRORQ, CROSS, CARRY, SIN, COS, (STH), (FIL), PLOT1, PLOT2, PLOT3, FPL0T4
Common Requirements	- None

-7-

```

*      LIST
*      LABEL
CFILTP  FILET WITH PLOT
        SUBROUTINE FILET (CP,CQ,CR,CS,XP,XQ,XR,XS,DUSE,CON,
1          THIC,ANS,N,L,FIX,SPACE,NREP)
90      FORMAT (1H , 10E12.5)
100     FORMAT(1H1/24H  TEMP VS TIME  AT DEPTH , F7.2,4H CM. /1H0 )
101     FORMAT(1H0,80X,28H TIME IN HOURS FROM SUNRISE  /1H0,
1        16H      TIME STEPS  ,F7.2,6H HOURS  )
22      FORMAT(1H1)
        DIMENSION N(25),CP(25),CQ(25),CR(25),CS(25),DUSE(5),L(5)
        DIMENSION ANS(999),XP(25),XQ(25),XR(25),XS(25),CON(5),THIC(5)
        DIMENSION TIME(999),ARRAY(800)
        NREP=NREP
        DO 1 I=1,NREP,1
        ANS(I)=0.0
1        CONTINUE
        K=0
2        M=0
        K=K+1
        IF (N(K)) 9,8,8
8        X=FIX
        TERM=N(K)
3        M=M+1
        XCAR=X
        IF (L(M)) 4,5,6
4        CALL ERRORQ (5,2)
5        IF (X-THIC(M)) 6,7,7
7        XCAR=THIC(M)
6        CALL CROSS (XP(K),XQ(K),XR(K),XS(K),CP(K),CQ(K),CR(K),CS(K),
1          DUSE(M),CON(M),EGA,ZETA,TERM)
        CALL CARRY (XP(K),XQ(K),XR(K),XS(K),CP(K),CQ(K),CR(K),CS(K),
1          CON(M),EGA,ZETA,XCAR)
        IF (L(M)) 4,13,11
13       IF (X-THIC(M)) 11,11,10
10       X=X-THIC(M)
        GO TO 3
11       TAU=0.0
        DO 12 I=1,NREP,1
        ANS(I)=ANS(I)+CP(K)*COSF(EGA*TAU)+CQ(K)*SINF(EGA*TAU)
        TIME(I)=TAU
        TAU=TAU+SPACE
12       CONTINUE
        GO TO 2
9        DO 14 I=1,NREP,10
        K=I+9
        WRITE OUTPUT TAPE 3,90,(ANS(J),J=I,K)
14       CONTINUE
        DO 202 I=1,NREP
        IF(ANS(I)-100.0)21,201,201
201      IF(ANS(I)-400.0)202,21,21
202      CONTINUE

```

-8-

```
20  NVL=((NREP-1)/10)+1
    NSBV=100/(NVL+1)
    CALL PLOT1(0,5,10,NVL,NSBV)
    CALL PLOT2(ARRAY,TIME(NREP)+SPACE,0.0,400.0,100.0)
    CALL PLOT3(1H*,TIME,ANS,NREP)
    WRITE OUTPUT TAPE 3,100,FIX
    CALL FLOT4(36,36H      TEMPERATURE  *DEGREES KELVIN*  )
    WRITE OUTPUT TAPE 3,101,SPACE
21  WRITE OUTPUT TAPE 3,22
    RETURN
    END
```

-E9-

**Entry Name** - FILEX

**Category** - Subroutine - Phase 1 only.

**Purpose** - To calculate and report as output a profile of normal temperatures vs. depth for fixed time.

**Arguments**

- CP, CQ, CR, CS, XP, XQ, XR, XS, DUSE, CON, THIC, ANS, N, L, FIX, SPACE, NREP
- L : The vector L(M) which indicates type of definition for thickness of each layer M,
- FIX : The fixed time for the profile,
- SPACE : Depth between successive values in the profile,
- NREP : Number of values in the profile beginning with value at  $x = 0$ .

All other arguments appear in the glossary.

**Error Exits** - None

**Description** - In the original concept, it seemed likely that this subroutine would be used. The necessary coding was postponed until needed. No need arose. Hence, the present version simply prints an appropriate message and returns.

**Lower Memory Requirements** - 30

**Transfer Vector** - (STH), (FIL)

**Common Requirements** - None

-10-

● LIST  
● LABEL  
SUBROUTINE FILEX (CP,CQ,CR,CS,XP,XQ,XR,XS,DUSE,CON,  
1 THIC,ANS,N,L,FIX,SPACE,NREP)  
90 FORMAT (1H ,31HX PROFILE SUBRTN. IS INCOMPLETE)  
DIMENSION N(25),CP(25),CQ(25),CR(25),CS(25),DUSE(5),L(5)  
DIMENSION ANS(999),XP(25),XQ(25),XR(25),XS(25),CON(5),THIC(5)  
WRITE OUTPUT TAPE 3,90  
RETURN  
END

-E11-

**Entry Name** - CROSS

**Category** - Subroutine - Phase 1 only.

**Purpose** - To apply equation A4 for a single value of  $\omega \neq 0$  (or A2a for  $\omega = 0$ ).

**Arguments** - XP, XQ, XR, XS, CP, CQ, CR, CS, DUSE, CON, EGA, ZETA, TERM

TERM : N(K) (floating point) to be used in calculating EGA.

All other arguments appear in the glossary.

**Error Exits** - #5 TERM is negative

**Description** - In addition to P, Q, R, and S, EGA and ZETA are also calculated.

**Lower Memory Requirements** - 122

**Transfer Vector** - SQRT, ERRORQ

**Common Requirements** - None

-12-

```
* LIST
* LABEL
CLUNARI ALPHABET FOR LABEL
SUBROUTINE CROSS (XP,XQ,XR,XS,CP,CQ,CR,CS,DUSE,CON,EGA,ZETA,TERM)
EGA=0.0088656*TERM
ZETA=SQRTF(0.5*EGA/DUSE)
IF (TERM) 1,2,3
1 CALL ERRORQ (5,2)
2 XP=CP
  XQ=CQ
  XR=-CR/CON
  GO TO 4
3 SP=0.5*CP
  SQ=0.5*CQ
  USE=-4.0*ZETA*CON
  SR=CR/USE
  SS=CS/USE
  XP=SP+SR-SS
  XQ=SQ+SR+SS
  XR=SP-SR+SS
  XS=SQ-SR-SS
4 CONTINUE
  RETURN
  END
```

-E13-

Entry Name - CARRY

Category - Subroutine - Phase 1 only.

Purpose - To apply equation A6 and A7 for a single value of  $\omega \neq 0$  (or A2a if  $\omega = 0$ ).

Arguments - XP, XQ, XR, XS, CP, CQ, CR, CS, CON, EGA, ZETA, X. All other arguments appear in the glossary.  
X : Depth at which A7 is to be evaluated (or A2a if  $\omega = 0$ )

Error Exits - #5 EGA is negative

Lower Memory Requirements - 153

Transfer Vector - ERRORQ, EXP, COS, SIN

Common Requirements - None



-14-

```
* LIST
* LABEL
CLUNARI ALPHABET FOR LABEL
SUBROUTINE CARRY (XP,XQ,XR,XS,CP,CQ,CR,CS,CON,EGA,ZETA,X)
IF (EGA) 1,2,3
1 CALL ERRORQ (5,2)
2 CP=XP+XR*X
  CQ=XQ
  GO TO 4
3 ZX=ZETA*X
  ER=EXPF(ZX)
  CE=COSF(ZX)
  SE=SINF(ZX)
  SP=(XP*CE+XQ*SE)*ER
  SQ=(XQ*CE-XP*SE)*ER
  SR=(XR*CE-XS*SE)/ER
  SS=(XR*SE+XS*CE)/ER
  CP=SP+SR
  CQ=SQ+SS
  CR=-CON*ZETA*(SP-SR+SQ-SS)
  CS=CON*ZETA*(SP-SR-SQ+SS)
4 CONTINUE
  RETURN
  END
```

-E15-

**Entry Name** - ERRORQ and DMPCOR

**Category** - Subroutines - Both Phase 1 and Phase 2.

**Purpose** - To report the nature of and dump core due to a serious error.

**Arguments** - ITYPE, IDUMP for ERRORQ.

ITYPE : Error identification number,  
IDUMP : Signal for DUMP,  
          = 1 return after printing ITYPE,  
          ≠ 1 call DMPCOR and DUMP.  
DMPCOR has no arguments.

**Description** - DMPCOR dumps all of core with mnemonics.

**Lower Memory Requirements** - 57 and 11

**Transfer Vector** - (SPH), (FIL), (STH), DMPCOR for ERRORQ  
DUMP for DMPCOR

**Common Requirements** - None

-16-

```
* LIST
* LABEL
SUBROUTINE ERRORQ (ITYPE,IDUMP)
92 FORMAT (1H1,10HERROR NO. ,I5)
93 FORMAT (1H ,12HDUMP FOLLOWS)
PRINT 92,ITYPE
WRITE OUTPUT TAPE 3,92,ITYPE
IF (IDUMP-1) 2,1,2
2 PRINT 93
WRITE OUTPUT TAPE 3,93
CALL DMPCOR
1 CONTINUE
RETURN
END
```

```
* LIST
* LABEL
* FAP
COUNT 2
*ERRORQ ALPHABET TO LABEL DMPCOR FOR INCLUSION IN BINARY DECK ERRORQ
ENTRY DMPCOR
DMPCOR CALL DUMP,00,9999,3,10000,19999,3,20000,32767,3
END
```

-E17-

**Entry Name** - LUNAR2

**Category** - Main Program - Phase 2

**Purpose** - Compute temperature perturbations due to the disc and direct output as specified by control cards.

**Arguments** - None (main program)

**Dimension and Common Considerations** - Standard Phase 2 COMMON and DIMENSION statements are used. In addition, S(J,K) is set S(254,103).

**Error Exits** - #98 - specified problem will not fit in the dimension restriction of S.  
#99 - 7090 failure.

**Unusual Cautions** - The arrangement of data within the matrix S is dependent on the input parameters. Restrictions on input are considered on Page E42.

**Description** - Calculation is made as described in Appendix B after input parameters have been read. Output is via subroutine LUAU2.

**Lower Memory Requirements** - 943

**Transfer Vector** - (FPT), FOURCO, LAYERS, (TSH), (RTN), (STH), (FIL), LUAU2, EXIT, VALUE, RADIO, BIQUAD, ERRORQ.

**Common Requirements** - Standard LUNAR2 26,618.

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\* LIST 8

• LABEL

CLUNAR2 ALPHABET FOR LABEL

```

COMMON N, DP, DQ, DR, DS, CP, CQ, CR, CS, XP, XQ, XR, XS, EGA, ZETA, NFORCO, COM
COMMON THIC, RHO, SPEH, DUSE, NLAYER, FIX, SPACE, DELT, DELX, DELR, TSXS
COMMON NDELT, NDP
COMMON TSRS, NFDS, NSDS, NXANS, NRANS, NDISC, ESH, TSH, VIS, Q, ESU, SIG, S
DIMENSION N(25), DP(25), DQ(25), DR(25), DS(25), EGA(25), ZETA(25)
DIMENSION CP(25), CQ(25), CR(25), CS(25), XP(25), XQ(25), XR(25), XS(25)
DIMENSION CON(10), THIC(10), RHO(10), SPEH(10), DUSE(10), NDP(10)
DIMENSION S(254, 103)
85  FORMAT (1H0, 15HINITIAL TIME = , 1E12.4)
89  FORMAT (6E12.4)
88  FORMAT (6I7)
92  FORMAT (1X, 6HDELT= , E10.4, 8H DELX= , E10.4, 8H DELR= , E10.4)
93  FORMAT (1X, 6HDELT= , I5, 8H NXANS= , I5, 8H NRANS= , I5, 8H NDISC= , I5,
1    8H NFDS= , I5, 7H NSDS= , I5, 10H HEIGHT= , E10.4)
    CALL FOURCO (2, 1)
    CALL LAYERS (1)
    READ INPUT TAPE 2, 89, ESH, TSH, VIS, Q, ESU, SIG
    READ INPUT TAPE 2, 89, DELT, DELX, DELR, TIN, HEIGHT
    READ INPUT TAPE 2, 88, NDELT, NXANS, NRANS, NDISC, NFDS, NSDS
    WRITE OUTPUT TAPE 3, 92, DELT, DELX, DELR
    WRITE OUTPUT TAPE 3, 85, TIN
    WRITE OUTPUT TAPE 3, 93, NDELT, NXANS, NRANS, NDISC, NFDS, NSDS, HEIGHT
    DO 18 I=1, NLAYER, 1
    NDP(I)=THIC(I)/DELX+0.5
18  CONTINUE
    NLAYER=NLAYER
    NDP(NLAYER)=32000
    A=DELX*ESU*SIG/CON(1)
    A=(2.0*A)/3.0
    FIX=0.0
    DXS=DELX*DELX
    DRS=DELR*DELR
    TSXS=DELT/DXS
    TSRS=DELT/DRS
    NDX=1
    MS=NSDS-1
    JR=NDISC-2
    IF (NDELT-NXANS) 31, 31, 32
31  NTA=NDELT
    NTB=0
    GO TO 76
32  NTA=(NXANS+NDELT)/2
    NTB=(NDELT-NXANS+1)/2
76  JINK=1
    IF (NDISC) 75, 75, 33
75  JINK=-1
    JR=NRANS+NDELT+1
    NDISC=10000
    IF (JR-NFDS) 37, 98, 98
33  JR=JR+1

```

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```

36 IF (JR+1+NTA-NTB-NRANS) 33,36,36
37 IF (JR+1+NTA-NFDS) 37,98,98
38 IF (NTA+2-NSDS) 38,38,98
38 CONTINUE
   DO 1 J=1,NFDS,1
   DO 39 K=1,NSDS,1
   S(J,K)=0.0
39 CONTINUE
1  CONTINUE
   MB=MS+2
   NTIM=0
   T=0.0
   DO 24 NT=1,NTA,1
   JR=JR+JINK
   T=T+DELT
   NTIM=NTIM+1
   TIME=TIN+T
   MB=MB-1
   IF (MS-MB) 17,22,22
22 GO TO 2
2  LAY=1
   NEDP=NDP(1)
   LEX=0
   DO 9 MX=MB,MS,1
   M=MX
   LEX=LEX+1
   IF (LEX-NEDP) 5,3,6
3  IF (MS-MX) 99,7,4
4  LAY=LAY+1
   M=M+1
5  MSTO=M-1
   GO TO 10
6  LAY=LAY-1
   M=M-1
7  MSTO=M-1
   GO TO 12
8  LAY=LAY+1
   NEDP=NEDP+NDP(LAY)
9  CONTINUE
17 M=MB-1
   MSTO=M-1
   GO TO 14
24 CONTINUE
   IF (NTB) 99,34,35
35 NDX=2
   MU=NTA
   MLAY=LAY
   NEDP=NEDP-NDP(LAY)
   DO 25 NT=1,NTB,1
   JR=JR-1
   T=T+DELT
   NTIM=NTIM+1
   TIME=TIN+T

```

-20-

```

NU=MU
MU=MU-1
LAY=MLAY
MB=NSDS-1
MEDP=NEDP
DO 20 MX=1, MU, 1
MB=MB-1
M=MB
MSTO=M+1
NU=NU-1
IF (NU-MEDP) 28, 26, 10
26 IF (MX-1) 99, 29, 27
27 LAY=LAY-1
M=M-1
MSTO=MSTO-1
GO TO 10
28 M=M+3
MSTO=MSTO+1
GO TO 12
29 LAY=LAY-1
MLAY=LAY
NEDP=NEDP-NDP(LAY)
DO 30 J=1, JR, 1
S(J, MSTO)=0.0
30 CONTINUE
23 MEDP=MEDP-NDP(LAY)
20 CONTINUE
M=MSTO
MSTO=MSTO-1
GO TO 14
25 CONTINUE
34 CONTINUE
CALL LUAU2 (NTIM, MSTO, JR)
CALL EXIT
16 CONTINUE
CALL LUAU2 (NTIM, MSTO, JR)
GO TO (24, 25), NDX
10 CNR=0.0
DO 11 J=2, JR, 1
CNR=CNR+2.0
ANS=S(J, M)+DUSE(LAY)*((S(J, M-1)+S(J, M+1)-2.0*S(J, M))*TSXS
1 +((S(J+1, M)-S(J-1, M))/CNR+S(J+1, M)+S(J-1, M)-2.0*S(J, M))*TSRS)
S(J, MSTO)=ANS
11 CONTINUE
S(1, MSTO)=S(2, MSTO)
GO TO (9, 20), NDX
12 DO 13 J=2, JR, 1
ANS=(CON(LAY)*S(J, M-2)+CON(LAY+1)*S(J, M))/(CON(LAY)+CON(LAY+1))
S(J, MSTO)=ANS
13 CONTINUE
S(1, MSTO)=S(2, MSTO)
GO TO (8, 23), NDX

```

-21-

```
14 CALL VALUE (0.0, TIME, V, F)
   R=0.0
   DO 15 J=2, JR, 1
   R=R+DELR
   RAD=HEIGHT
   CALL RADIO (R, TIME, RAD)
   B=V + ( (4.0*S(J, M) - S(J, M+1) + 2.0*DELR*(RAD+F))/3.0 )
   CALL BIQUAD(A, B, 400.0, 0.001, ANS)
   S(J, MSTO)=ANS-V
15 CONTINUE
   S(1, MSTO)=S(2, MSTO)
   GO TO 16
99 CALL ERRORQ (99, 2)
98 CALL ERRORQ (98, 2)
END
```



-E22-

Entry Name	- LUNE2M
Category	- Main Program - Phase 2.
Purpose	- Compute temperature perturbations due to the disc and direct output as specified by control cards.
Arguments	- None (main program)
Dimension and Common Considerations	- Standard Phase 2 COMMON and DIMENSION statements are used. In addition, S(K) is set S(25000).
Error Exits	- #98 - specified problem will not fit in the dimension restriction of S. #99 - 7090 failure.
Unusual Cautions	- The size of S(K) needed is dependent on the input parameters. Restrictions on input are considered on Page E42.
Description	- Calculation is made as described in Appendix B after input parameters have been read assuming that the independent variable, r(radial distance) is irrelevant. Output is via subroutine LUNO2.
Lower Memory Requirements	- 775
Transfer Vector	- (FPT), FOURCO, LAYERS, (TSH), (RTN), (STH), (FIL), LUNO2, VALUE, RADIO, BIQUAD, ERRORQ.
Common Requirements	- Standard LUNE2M 25,456.

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```

* LIST 8
* LABEL
CLUNE2M LUNAR2 MODIFIED FOR ONE-D. PROBLEM
COMMON N, DP, DQ, DR, DS, CP, CQ, CR, CS, XP, XQ, XR, XS, EGA, ZETA, NFORCO, CON
COMMON THIC, RHO, SPEH, DUSE, NLayer, FIX, SPACE, DELT, DELX, DELR, TSXS
COMMON NDELT, NDP
COMMON TPRS, NFDS, NSDS, NXANS, NRANS, NDISC, ESH, TSH, VIS, Q, ESU, SIG, S
DIMENSION N(25), DP(25), DQ(25), DR(25), DS(25), EGA(25), ZETA(25)
DIMENSION CP(25), CQ(25), CR(25), CS(25), XP(25), XQ(25), XR(25), XS(25)
DIMENSION CON(10), THIC(10), RHO(10), SPEH(10), DUSE(10), NDP(10)
DIMENSION S(25000)
85 FORMAT (1H0,15HINITIAL TIME = ,1E12.4)
89 FORMAT (6E12.4)
88 FORMAT (6I7)
92 FORMAT (1X,6HDELT= ,E10.4,8H DELX= ,E10.4,8H DELR= ,E10.4)
93 FORMAT (1X,6HNDDEL= ,I5,8H NXANS= ,I5,8H NRANS= ,I5,8H NDISC= ,I5,
1 8H NFDS= ,I5,7H NSDS= ,I5,10H HEIGHT= ,E10.4)
100 FORMAT (1X,1A8)
CALL FOURCO (2,1)
102 CALL LAYERS (1)
READ INPUT TAPE 2,89,ESH,TSH,VIS,Q,ESU,SIG
READ INPUT TAPE 2,89,DELT,DELX,DELR,TIN,HEIGHT
READ INPUT TAPE 2,88,NDELT,NXANS,NRANS,NDISC,NFDS,NSDS
WRITE OUTPUT TAPE 3,92,DELT,DELX,DELR
WRITE OUTPUT TAPE 3,85,TIN
WRITE OUTPUT TAPE 3,93,NDELT,NXANS,NRANS,NDISC,NFDS,NSDS,HEIGHT
DO 18 I=1,NLayer,1
NDP(I)=THIC(I)/DELX+0.5
18 CONTINUE
NLayer=NLayer
NDP(NLayer)=32000
A=DELX*ESU*SIG/CON(1)
A=(2.0*A)/3.0
FIX=0.0
DXS=DELX*DELX
TSXS=DELT/DXS
NDX=1
NSDS=25000
MS=NSDS-1
NDISC=10000
JR=2
IF (NDELT-NXANS) 31,31,32
31 NTA=NDELT
NTB=0
GO TO 37

```

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```

32  NTA=(NXANS+NDELT)/2
    NTB=(NDELT-NXANS+1)/2
37  IF (NTA+2-NSDS) 38,38,98
38  CONTINUE
    DO 39 K=1,NSDS,1
    S(K)=0.0
39  CONTINUE
    MB=MS+2
    NTIM=0
    T=0.0
    DO 24 NT=1,NTA,1
    T=T+DELT
    NTIM=NTIM+1
    TIME=TIN+T
    MB=MB-1
    IF (MS-MB) 17,22,22
22  GO TO 2
2   LAY=1
    NEDP=NDP(1)
    LEX=0
    DO 9 MX=MB,MS,1
    M=MX
    LEX=LEX+1
    IF (LEX-NEDP) 5,3,6
3   IF (MS-MX) 99,7,4
4   LAY=LAY+1
    M=M+1
5   MSTO=M-1
    GO TO 10
6   LAY=LAY-1
    M=M-1
7   MSTO=M-1
    GO TO 12
8   LAY=LAY+1
    NEDP=NEDP+NDP(LAY)
9   CONTINUE
17  M=MB-1
    MSTO=M-1
    GO TO 14
24  CONTINUE
    IF (NTB) 99,34,35
35  NDX=2
    MU=NTA
    MLAY=LAY
    NEDP=NEDP-NDP(LAY)
    DO 25 NT=1,NTB,1
    T=T+DELT
    NTIM=NTIM+1
    TIME=TIN+T
    NU=MU
    MU=MU-1
    LAY=MLAY
    MB=NSDS-1

```

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```

MEDP=NEDP
DO 20 MX=1, MU, 1
MB=MB-1
M=MB
MSTO=M+1
NU=NU-1
IF (NU-MEDP) 28,26,10
26 IF (MX-1) 99,29,27
27 LAY=LAY-1
M=M-1
MSTO=MSTO-1
GO TO 10
28 M=M+3
MSTO=MSTO+1
GO TO 12
29 LAY=LAY-1
MLAY=LAY
NEDP=NEDP-NDP(LAY)
S(MSTO)=0.0
23 MEDP=MEDP-NDP(LAY)
20 CONTINUE
M=MSTO
MSTO=MSTO-1
GO TO 14
25 CONTINUE
34 CONTINUE
CALL LUNO2 (NTIM, MSTO, JR)
101 READ INPUT TAPE 2, 100, TEST
IF (TEST-6HREPEAT) 101, 102, 101
16 CONTINUE
CALL LUNO2 (NTIM, MSTO, JR)
GO TO (24, 25), NDX
10 CONTINUE
ANS=S( M)+DUSE(LAY)*((S( M-1)+S( M+1)-2.0*S( M))*TSXS )
S(MSTO)=ANS
GO TO (9, 20), NDX
12 CONTINUE
ANS=(CON(LAY)*S( M-2)+CON(LAY+1)*S( M))/(CON(LAY)+CON(LAY+1))
S(MSTO)=ANS
GO TO (8, 23), NDX
14 CALL VALUE (0.0, TIME, V, F)
R=0.0
RAD=HEIGHT
CALL RADIO (R, TIME, RAD)
B=V + ( (4.0*S( M)-S( M+1)+2.0*DELX*(RAD+F))/3.0 )
CALL BIQUAD(A, B, 400.0, 0.001, ANS)
S( MSTO)=ANS-V
GO TO 16
99 CALL ERRORQ (99, 2)
98 CALL ERRORQ (98, 2)
END

```

-E26-

Entry Name - LUAU2

Category - Subroutine - Phase 2.

Purpose - Output Subroutine for Phase 2.

Arguments - NTIM, MSTO, JR

NTIM : Number of time steps for which the Phase 2 calculation has been completed,

MSTO : Value of the column (last dimension of S) in which surface temperature perturbations currently appear,

JR : Number of r-steps to last non-zero perturbation.

Dimension and Common Considerations - Standard Phase 2 COMMON and DIMENSION statements are used. In addition, S(J,K) is set S(254,103).

Error Exits - None

Unusual Cautions - The dimensionality of S must agree with that given in the main program, LUNAR2.

Description - Output consists of radial profiles for fixed depth as specified by control cards read by LUAU2. If NTIM is sufficiently small, no output occurs.

Lower Memory Requirements - 223

Transfer Vector - (TSH), (RTN), (STH), (FIL)

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\*  
\*  
LIST

## LABEL

SUBROUTINE LUAU2 (NTIM, MSTO, JR)

COMMON N, DP, DQ, DR, DS, CP, CQ, CR, CS, XP, XQ, XR, XS, EGA, ZETA, NFORCO, COM

COMMON THIC, RHO, SPEH, DUSE, NLayer, FIX, SPACE, DELT, DELX, DELR, TSXS

COMMON NDELT, NDP

COMMON TSRS, NFDS, NSDS, NXANS, NRANS, NDISC, ESH, TSH, VIS, Q, ESU, SIG, S

DIMENSION N(25), DP(25), DQ(25), DR(25), DS(25), EGA(25), ZETA(25)

DIMENSION CP(25), CQ(25), CR(25), CS(25), XP(25), XQ(25), XR(25), XS(25)

DIMENSION CON(10), THIC(10), RHO(10), SPEH(10), DUSE(10), NDP(10)

DIMENSION S(254, 103)

DIMENSION ANS(10)

86 FORMAT (1H0, 19HRADIAL PROFILE. T= , E14.8, 5H X= , E14.8,

1 11H R-SPACE= , E14.8)

87 FORMAT (1E12.4, 3I7)

90 FORMAT (1H , 10E11.4)

IF (FIX) 2, 1, 2

1 READ INPUT TAPE 2, 87, FIX, MOX, MDRS, NDRS

NFI=FIX/DELT

2 IF (NFI-NTIM) 1, 3, 4

4 RETURN

3 LOX=MSTO+MOX

WORKA=MOX

WORKA=WORKA\*DELX

WORKB=MDRS

WORKB=WORKB\*DELR

WRITE OUTPUT TAPE 3, 86, FIX, WORKA, WORKB

K=1

10 DO 5 I=1, 10, 1

ANS(I)=0.0

5 CONTINUE

DO 9 J=1, 10, 1

IF (NSDS-1-LOX) 6, 7, 7

7 IF (JR-K) 6, 8, 8

8 ANS(J)=S(K, LOX)

6 K=K+MDRS

9 CONTINUE

WRITE OUTPUT TAPE 3, 90, (ANS(I), I=1, 10)

NDRS=NDRS-10

IF (NDRS) 1, 1, 10

END

-E28-

**Entry Name** - LUNO2

**Category** - Subroutine - Phase 2

**Purpose** - Output subroutine for Phase 2 modified for one-dimensional problem.

**Arguments** - NTIM, MSTO, JR

NTIM : Number of time steps for which the Phase 2 calculation has been completed,

MSTO : Value of the column (only dimension of S) in which surface temperature perturbations currently appear,

JR : Meaningless (LUNO2 is a modification of LUAU2).

**Dimension and Common Considerations** - Standard Phase 2 COMMON and DIMENSION statements are used. In addition, S(K) is set S(25000).

**Error Exits** - None

**Unusual Cautions** - The dimensionality of S must agree with that given in the main program, LUNE2M.

**Description** - Output consists of depth profiles for fixed time as specified by control cards read by LUNO2. If NTIM is sufficiently small, no output occurs.

**Lower Memory Requirements** - 188

**Transfer Vector** - (TSH), (RTN), (STH), (FIL)

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## \* LIST

## \* LABEL

```

SUBROUTINE LUNO2 (NTIM,MSTO,JR)
COMMON N,DP,DQ,DR,DS,CP,CQ,CR,CS,XP,XQ,XR,XS,EGA,ZETA,NFORCO,CON
COMMON THIC,RHO,SPEH,DUSE,NLAYER,FIX,SPACE,DELT,DELX,DELR,TSXS
COMMON NDELT,NDP
COMMON TSRS,NFDS,NSDS,NXANS,NRANS,NDISC,ESH,TSH,VIS,Q,ESU,SIG,S
DIMENSION N(25),DP(25),DQ(25),DR(25),DS(25),EGA(25),ZETA(25)
DIMENSION CP(25),CQ(25),CR(25),CS(25),XP(25),XQ(25),XR(25),XS(25)
DIMENSION CON(10),THIC(10),RHO(10),SPEH(10),DUSE(10),NDP(10)
DIMENSION S(25000)
DIMENSION ANS(10)
86  FORMAT (1H0,18HDEPTH PROFILE. T= ,E14.8,11H FIRST-X= ,E14.8,
1    11H X-SPACE= ,E14.8)
87  FORMAT (1E12.4,3I7)
90  FORMAT (1H ,10E11.4)
    IF (FIX) 2,1,2
1    READ INPUT TAPE 2,87,FIX,MOX,MDRS,NDRS
    NFIT=FIX/DELT
2    IF (NFIT-NTIM) 1,3,4
4    RETURN
3    LOX=MSTO+MOX
    WORKA=MOX
    WORKA=WORKA*DELX
    WORKB=MDRS
    WORKB=WORKB*DELX
    WRITE OUTPUT TAPE 3,86,FIX,WORKA,WORKB
10   DO 5 I=1,10,1
    ANS(I)=0.0
5    CONTINUE
    DO 9 J=1,10,1
    IF (NSDS-1-LOX) 8,7,7
7    ANS(J)=S(LOX)
8    LOX=LOX+MDRS
9    CONTINUE
    WRITE OUTPUT TAPE 3,90,(ANS(I),I=1,10)
    NDRS=NDRS-10
    IF (NDRS) 1,1,10
END

```



-E30-

Entry Name	- RADIO
Category	- Subroutine - Phase 2
Purpose	- Computes surface radiation input divided by conductivity of surface layer.
Arguments	- R, T, RAD  R : radius (cm. from center of disc), T : time (hours from nearest previous sunrise), RAD : input argument - height of disc above surface; output - surface radiation/CON(1).
Dimension and Common Considerations	- Standard Phase 2 COMMON and DIMENSION statements are used.
Error Exits	- None
Unusual Cautions	- T must be positive and less than one lunar cycle.
Description	- The angular effect of the geometry of the disc is calculated if height is positive, otherwise this effect is ignored.
Lower Memory Require- ments	- 179
Transfer Vector	- SIN, SQRT

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• LIST  
\* LABEL

CRADIOX RADIATION INPUT FOR LUNAR2 WITH RADIATION EDGE EFFECTS

SUBROUTINE RADIO(R,T,RAD)

COMMON N,DP,DQ,DR,DS,CP,CQ,CR,CS,XP,XQ,XR,XS,EGA,ZETA,NFORCO,CON

COMMON THIC,RHO,SPEH,DUSE,NLAYER,FIX,SPACE,DELT,DELX,DELR,TSXS

COMMON NDELT,NDP

COMMON TSRS,NFDS,NSDS,NXANS,NRANS,NDISC,ESH,TSH,VIS,Q,ESU,SIG,S

DIMENSION N(25),DP(25),DQ(25),DR(25),DS(25),EGA(25),ZETA(25)

DIMENSION CP(25),CQ(25),CR(25),CS(25),XP(25),XQ(25),XR(25),XS(25)

DIMENSION CON(10),THIC(10),RHO(10),SPEH(10),DUSE(10),NDP(10)

HEIGHT=RAD

ANDISC=NDISC

RSH=ANDISC\*DELR

IF(HEIGHT)8,1,7

1 RAD=0.0

IF(RSH-R)3,2,2

2 RAD=ESH\*ESU\*SIG\*TSH\*\*4/CON(1)

GO TO 8

3 IF(EGA(2)\*T-3.1415927)4,5,5

4 SOL=VIS\*Q\*SINF(EGA(2)\*T)/CON(1)

GO TO 6

5 SOL=0.0

6 RAD=RAD+SOL

GO TO 8

7 B=RSH/HEIGHT

C=R/HEIGHT

F=0.5(1.0-(1.0 + C\*\*2 - B\*\*2)/SQRTF(C\*\*4 + 2.0\*C\*\*2 \*(1.0-B\*\*2)

1 + (1.0 + B\*\*2)\*\*2))

RAD=ESH\*ESU\*SIG\*F\*TSH\*\*4/CON(1)

IF(RSH-R)3,5,5

8 RETURN

END

-E32-

Entry Name - LAYERS

Category - Subroutine - Phase 2.

Purpose - To read description of layers.

Arguments - IOUT  
           IOUT : Specifies output as follows:  
                    $\geq 1$  output of layer descriptions,  
                    $< 1$  no output layer descriptions.

Dimension and Common Considerations - Standard Phase 2 COMMON and DIMENSION Statements are used.

Error Exits - None

Unusual Cautions - L(M) is the fixed point variable for layer M which indicates:  
                   L(M) = 0 : layer thickness is defined,  
                   L(M) > 0 : layer thickness is undefined ( $\infty$ ),  
                   L(M) < 0 : Call EXIT immediately.  
                   In addition to reading L(M), CON(M), THIC(M), RHO(M), and SPEH(M), LAYERS computes DUSE(M) = CON(M)/(RHO(M)\* SPEH(M)).  
                   The infinite layer is, of course, read last and serves as a signal that all layering descriptions have then been read.  
                   LAYERS also sets NLAYER.

Lower Memory Requirements - 138

Transfer Vector - (STH), (FIL), (TSH), (RTN), EXIT.

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```

* LIST
* LABEL
SUBROUTINE LAYERS (IOUT)
90  FORMAT (1X,1I3,4E14.8)
94  FORMAT (1H0,18HPHYSICAL CONSTANTS)
95  FORMAT (1H,1I3,4E14.8)
COMMON N, DP, DQ, DR, DS, CP, CQ, CR, CS, XP, XQ, XR, XS, EGA, ZETA, NFORCO, CON
COMMON THIC, RHO, SPEH, DUSE, NLAYER, FIX, SPACE, DELT, DELX, DELR, TSXS
COMMON NDELT, NDP
COMMON TSRS, NFDS, NSDS, NXANS, NRANS, NDISC, ESH, TSH, VIS, Q, ESU, SIG, S
DIMENSION N(25), DP(25), DQ(25), DR(25), DS(25), EGA(25), ZETA(25)
DIMENSION CP(25), CQ(25), CR(25), CS(25), XP(25), XQ(25), XR(25), XS(25)
DIMENSION CON(10), THIC(10), RHO(10), SPEH(10), DUSE(10), NDP(10)
K=0
IF (IOUT) 2,2,1
1  WRITE OUTPUT TAPE 3, 94
2  K=K+1
   READ INPUT TAPE 2, 90, L, CON(K), THIC(K), RHO(K), SPEH(K)
   IF (IOUT) 4,4,3
3  WRITE OUTPUT TAPE 3, 95, L, CON(K), THIC(K), RHO(K), SPEH(K)
4  IF (L) 7,5,8
8  THIC(K)=1.0E30
5  DUSE(K)=CON(K)/(RHO(K)*SPEH(K))
   IF (L) 7,2,6
6  NLAYER=K
   RETURN
7  CALL EXIT
END

```

-E34-

**Entry Name** - FOURCO

**Category** - Subroutine - Phase 2.

**Purpose** - Read Lunar Fourier coefficients and prepare these coefficients for further use.

**Arguments** - IPROB, IOUT

IPROB : specifies output title as follows:  
     $\leq 1$  title for Phase 1,  
     $= 2$  title for Phase 2,  
     $> 2$  no title.

IOUT : specifies output of Fourier coefficients:  
     $\leq 0$  no output of coefficients,  
     $> 0$  output of coefficients for reference.

**Dimension and Common Considerations** - Standard Phase 2 COMMON and DIMENSION statements are used.

**Error Exits** - None

**Description** - In addition to reading of Fourier coefficients into DP(K), DQ(K), DR(K), and DS(K); FOURCO reads N(K), calculates EGA(K), rescales DR(K) and DS(K) for time units in hours, and sets NFORCO.

**Lower Memory Requirements** - 150

**Transfer Vector** - (STH), (FIL), (TSH), (RTN).

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```

* LIST
* LABEL
SUBROUTINE FOURCO (IPROB,IOUT)
90 FORMAT (1X,1I3,4E14.8)
91 FORMAT (1H0,20HFOURIER COEFFICIENTS)
93 FORMAT (1H1,27HLUNAR1 TEMPERATURE PROFILES)
95 FORMAT (1H ,1I3,4E14.8)
98 FORMAT (1H1,32HLUNAR2 TEMPERATURE PERTURBATIONS)
COMMON N,DP,DQ,DR,DS,CP,CQ,CR,CS,XP,XQ,XR,XS,EGA,ZETA,NFORCO,CON
COMMON THIC,RHO,SPEH,DUSE,NLAYER,FIX,SPACE,DELT,DELX,DELR,TSXS
COMMON NDELT,NDP
COMMON TSRS,NFDS,NSDS,NXANS,NRANS,NDISC,ESH,TSH,VIS,Q,ESU,SIG,S
DIMENSION N(25),DP(25),DQ(25),DR(25),DS(25),EGA(25),ZETA(25)
DIMENSION CP(25),CQ(25),CR(25),CS(25),XP(25),XQ(25),XR(25),XS(25)
DIMENSION CON(10),THIC(10),RHO(10),SPEH(10),DUSE(10),NDP(10)
K=0
IF (IPROB-2) 1,2,3
1 WRITE OUTPUT TAPE 3,93
GO TO 3
2 WRITE OUTPUT TAPE 3,98
3 IF (IOUT) 5,5,4
4 WRITE OUTPUT TAPE 3,91
5 K=K+1
READ INPUT TAPE 2,90,N(K),DQ(K),DP(K),DS(K),DR(K)
DS(K)=3600.0*DS(K)
DR(K)=3600.0*DR(K)
TERM=N(K)
EGA(K)=0.0088656*TERM
IF (IOUT) 7,7,6
6 WRITE OUTPUT TAPE 3,95,N(K),DQ(K),DP(K),DS(K),DR(K)
7 IF (N(K)) 8,5,5
8 NFORCO=K-1
RETURN
END

```

-E36-

Entry Name - VALUE

Category - Subroutine - Phase 2

Purpose - To calculate temperature and x-derivative of temperature.

Arguments - X, T, V, F.  
X : Depth at which results are desired.  
T : Time (measured from nearest previous sunrise) at which results are desired.  
V : Temperature.  
F : x-derivative of temperature.

Dimension and Common Considerations - Standard Phase 2 COMMON and DIMENSION statements are used.

Error Exits - #10 - the argument, X, is negative. Also, the argument, X, is positive.

Unusual Cautions - In the original concept, it seemed likely that this subroutine would be used for positive X. The necessary coding was postponed until needed. No need arose.

Description - Using the Fourier coefficients in COMMON, Phase 1 V and F are calculated.

Lower Memory Requirements - 88

Transfer Vector - ERRORQ, COS, SIN

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```

• LIST
* LABEL
SUBROUTINE VALUE (X,T,V,F)
COMMON N,DP,DQ,DR,DS,CP,CQ,CR,CS,XP,XQ,XR,XS,EGA,ZETA,NFORCO,CON
COMMON THIC,RHO,SPEH,DUSE,NLAYER,FIX,SPACE,DELT,DELX,DELR,TSXS
COMMON NDELT,NDP
COMMON TSRS,NFDS,NSDS,NXANS,NRANS,NDISC,ESH,TSH,VIS,Q,ESU,SIG,S
DIMENSION N(25),DP(25),DQ(25),DR(25),DS(25),EGA(25),ZETA(25)
DIMENSION CP(25),CQ(25),CR(25),CS(25),XP(25),XQ(25),XR(25),XS(25)
DIMENSION CON(10),THIC(10),RHO(10),SPEH(10),DUSE(10),NDP(10)
IF (X) 3,4,5
5 CONTINUE
3 CALL ERRORQ (10,2)
4 V=0.0
  F=0.0
  DO 1 J=1,NFORCO,1
    TAU=T*EGA(J)
    CUS=COSF(TAU)
    SAN=SINF(TAU)
    V=V+DP(J)*CUS+DQ(J)*SAN
    F=F+DR(J)*CUS+DS(J)*SAN
1 CONTINUE
2 F=-F/CON(1)
  RETURN
  END

```



-E38-

Entry Name - BIQUAD

Category - Subroutine - Phase 2.

Purpose - To find the positive root of  $g(U) = AU^4 + U - B$  where  $A, B > 0$ .

Arguments - A, B, UMAX, FRACT, ROOT  
           A, B : Coefficients of  $g(0)$ .  
           UMAX : Maximum allowable root ( $^{\circ}K$ ,  $\therefore$  positive); a root is assumed to lie in  $(0, UMAX)$ .  
           FRACT: A root is supposed found if known to lie in a U-interval of length not exceeding FRACT.  
           ROOT : The desired positive root.

Error Exits - #6 - no root discovered in  $(0, UMAX)$ .

Unusual Cautions - If A and B are such that the root would lie outside  $(0, UMAX)$ , strange results occur. We successfully used,  $UMAX = 400.0$ .

Description - The root search uses an alternating half interval - secant intersection method.

Lower Memory Requirements - 141

Transfer Vector - ERRORQ

Common Requirements - None

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```
* LIST
* LABEL
SUBROUTINE BIQUAD (A, B, UMAX, FRACT, ROOT)
C SUBROUTINE BIQUAD, ROOT OF BIQUADRATIC EQUATION, LUNAR2
  QUARTF(X)=(A*X**4)+X-B
  X1=0.0
  X2=UMAX
  DO 7 INDEX=1, 100
    Y1=QUARTF(X1)
    Y2=QUARTF(X2)
    XPRIME=((X1*Y2)-(Y1*X2))/(Y2-Y1)
    YPRIME=QUARTF(XPRIME)
    IF(YPRIME) 2, 1, 1
1  ROOT=XPRIME
    GO TO 8
2  X1=XPRIME
    XPRIME=(X1+X2)/2.0
    YPRIME=QUARTF(XPRIME)
    IF(YPRIME) 3, 1, 5
3  X1=XPRIME
    GO TO 4
5  X2=XPRIME
4  IF(X2-X1-FRACT) 6, 6, 7
6  ROOT=(X1+X2)/2.0
    GO TO 8
7  CONTINUE
    CALL ERRORQ (6,2)
8  RETURN
  END
```

GLOSSARY OF NAMESI. Data Read by Both Phases

N(K) : Multipliers of  $\Omega$  in calculating  $\alpha_n$ .  
 DP(K) : Temperature cosine coefficients.  
 DQ(K) : Temperature sine coefficients.  
 DR(K) : Flux cosine coefficients.  
 DS(K) : Flux sine coefficients.  
 CON(M) : Layer thermal conductivities (cals/hr cm  $^{\circ}$ K).  
 THIC(M) : Layer thicknesses (cm).  
 RH $\phi$ (M) : Layer densities (gm/cm<sup>3</sup>).  
 SPEH(M) : Layer specific heats (cal/gm  $^{\circ}$ K).

II. Data Read by Phase 2 Only

ESH : Emissivity of disc (we used 0.9).  
 TSH : Temperature of disc (we used 300 $^{\circ}$ K).  
 VIS : Average reflectivity,  $\alpha_v$  (we used 0.875).  
 Q : Solar constant,  $q_0$  (we used 117.0 cal/cm<sup>2</sup> hr).  
 ESU : Emissivity of lunar surface,  $\epsilon_s$  (we used 0.9).  
 SIG : Stefan-Boltzmann constant,  $\sigma$  (cal/cm<sup>2</sup> hr ( $^{\circ}$ K)<sup>4</sup>).  
 DELT :  $\Delta t$  = time step (cm).  
 DELX :  $\Delta x$  = depth step (cm).  
 DELR :  $\Delta r$  = radial step (cm).  
 TIN : Initial time measured from first sunrise prior to placement of disc (hr).  
 HEIGHT : Height of disc above surface (cm).  
 NDEL T : Number of time steps to be taken in this run.  
 NXANS : Number of depth steps for which answers are desired after NDEL T time steps.  
 NRANS : Number of radial steps for which answers are desired after NDEL T time steps.  
 NDISC : Number of radial steps from center of disc to nearest grid point outside the radius of the disc.

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NFDS : First dimension of S.  
 NSDS : Second dimension of S.

### III. Variables Calculated by Both Phases

DUSE(M) : Layer thermal diffusivities.

### IV. Work Areas - Phase 1

CP(K), CQ(K), CR(K), and CS(K): Storage for data Fourier coefficients and results related to Equations A5 and A6.

XP(K), XQ(K), XR(K), and XS(K): Results of Equation A4.

EGA : Storage for  $\omega_n$  of current interest.

ZETA : Storage for  $\zeta_n$  of current interest.

### V. Work Areas - Phase 2

EGA(K) : Storage for all  $\omega_n$ .

### VI. Variables Calculated by Phase 2 Only

NFORCO : Number of values of  $\omega_n$  (including  $\omega_0 = 0$ ).

NLAYER : Number of layers.

TSXS :  $\Delta t / (\Delta x)^2$ .

TSRS :  $\Delta t / (\Delta r)^2$ .

NDP(M) : Number of depth steps required for layer M.  
 Defined to be 32,000 for an infinite layer.

S(J,K) : Storage area and work area for results in LUNAR2.

S(K) : Storage area and work area for results in LUNO2  
 (the one-dimensional modification of LUNAR2).

INPUT SEQUENCE AND RESTRICTIONSPhase 1:

## 1. Fourier Coefficients.

90     FORMAT (1X, 1I3, 4E14.8)

READ INPUT TAPE 2, 90, N(K), DQ(K), DP(K), DS(K),  
DR(K)

Iterate: K = 0,1,2, ....., M - 1, M

with    N(K) = 0,1,2, ....., M - 1, -1

N(M) = -01 signals termination of read.

DQ(M), DP(M), DS(M), DR(M) are arbitrary.

## 2. Layer Definition.

READ INPUT TAPE 2, 90, L(K), CON(K), THIC(K),  
RHO(K), SPEH(K)

Iterate: K = 0,1,2, ....., M - 1, M

L(K) = 0,0,0, ....., 0, 1

L(M) = 1 signals termination of read.

Layer M is infinite in thickness.

L(M) = -1 signals immediate exit.

## 3. Profile Specification.

92     FORMAT (1X, 1I3, 2E14.8, 1I4)

READ INPUT TAPE 2, 92, KFIL, FIX, SPACE, NREP

KFIL > 0 → Depth profile is desired,

KFIL = 0 → Time profile is desired,

KFIL < 0 → Return to 2 for new media.

See pages E4 and E9 for further information.

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Phase 2 (LUNAR2: S(J,K)):

1. Fourier Coefficients - Precisely as in Phase 1.
2. Layer Definition - Precisely as in Phase 1.
3. Parameters and Specifications.

89      FORMAT (6E12.4)

88      FORMAT (6I7)

READ INPUT TAPE 2, 89, ESH, TSH, VIS, Q, ESU, SIG

READ INPUT TAPE 2, 89, DELT, DELX, DELR, TIN, HEIGHT

READ INPUT TAPE 2, 88, NDELT, NXANS, NRANS, NDISC,  
NFDS, NSDS

Restrictions.

- a)  $DELR = DELX$  and  $DELT$  satisfy Eqn. B7 for the largest diffusivity in the problem.  $THIC(K)/DELX$  an integer for each  $K$  with finite  $THIC(K)$ .
- b) NFDS and NSDS are the first and second numbers in the DIMENSION statement defining  $S(J,K)$ .
- c) Denote "greatest integer  $\leq \alpha$ " by  $[\alpha]$ , and define

$$NTA = \text{MIN} ([ (NXANS + NDELT)/2 ], NDELT),$$

$$NTB = \text{MAX} ([ (NDELT - NXANS + 1)/2 ], 0),$$

$$JR = \text{MAX} (NDISC - 1, NRANS - NTA - 1 + NTB).$$

Then it must be true that

$$JR + 1 + NTA < NFDS, \text{ and } 2 + NTA \leq NSDS$$

where  $NDISC = [(15.24 \text{ cm}/DELR) + 0.99]$ .

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If these conditions cannot be satisfied, it will be necessary to re-dimension S.

It should be noted that LUNAR2 is specialized if the NDISC definition above is violated to the extent that NIDSC is not positive. This results in an inefficient use of LUNAR2 to solve the one-dimensional problem. This option should not be used since LUNE2M does the same work in less time.

#### 4. Radial Profile Specification (Read by LUAU2, P. E26).

87      FORMAT (1E12.4, 3I7)

READ INPUT TAPE 2, 87, FIX, MOX, MDRS, NDRS

FIX:      Time at which a radial profile is desired. FIX should increase monotonically from profile to profile. FIX/DELT should be an integer.

MOX:      Number of depth steps at which the profile is desired.

MDRS:      Number of radial steps between values on the profile beginning with the  $r = 0$  value.

NDRS:      Number of values in the profile.

#### Phase 2 (LUNE2M: S(K)):

1. Fourier Coefficients - Precisely as above.
2. Layer Definition - Precisely as above.
3. Parameters and Specifications - Read Precisely as above.

#### Restrictions.

- a) DELX and DELT satisfy Eqn. B7 for the largest

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-E45-

diffusivity of the problem.  $THIC(K)/DELX$  an integer for each  $K$  with finite  $THIC(K)$ .  $DELR$ ,  $NRANS$ ,  $NDISC$ ,  $NFDS$ , and  $NSDS$  are arbitrary.

- b) Denote "greatest integer  $\leq \alpha$ " by  $[\alpha]$ , and define.

$$NTA = \text{MIN} ([ (NXANS + NDELT)/2 ], NDELT).$$

Then it must be true that

$$NTA + 2 \leq 25,000.$$

If not, the problem cannot be solved.

#### 4. Depth Profile Specification (Read by LUNO2, P. E28).

87      FORMAT (1E12.4, 3I7)

READ INPUT TAPE 2, 87, FIX, MOX, MDRS, NDRS

**FIX:** Time at which a depth profile is desired. **FIX** should increase monotonically from profile to profile.  $FIX/DELT$  should be an integer.

**MOX:** Number of depth steps to the first (nearest surface) value in the profile.

**MDRS:** Number of depth steps between successive profile values.

**NDRS:** Number of values in the profile.

#### 5. Repetition Signal.

100      FORMAT (1X, 1A6)

READ INPUT TAPE 2, 100, TEST

If the word REPEAT is read into TEST return is made to 2 above for a new media.

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-E46-

Restriction.

In order for the word REPEAT to be properly interpreted it is necessary that the list of depth profiles described above (4) contain at least one profile whose time (FIX) is greater than the time obtained by an NDELT summation of DELT.